

Symplectic Morse Theory and Witten Deformation

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Abstract

On symplectic manifolds, we introduce a Morse-type complex with elements generated by pairs of critical points of a Morse function. The differential of the complex consists of gradient flows and an integration of the symplectic structure over spaces of gradient flow lines. Using the Witten deformation method, we prove that the cohomology of this complex is independent of both the Riemannian metric and the Morse function used to define the complex and is in fact isomorphic to the cohomology of differential forms of Tsai, Tseng and Yau (TTY). We also obtain Morse-type inequalities that bound the dimensions of the TTY cohomologies by the number of Morse critical points and the interaction of symplectic structure with the critical points.

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1 Introduction

The Morse complex, also referred to as the Morse-Witten or Smale-Thom complex, captures the information of the standard homology groups of a closed manifold M by means of a Morse function f , i.e. a function whose Hessian at each critical point is non-degenerate, and a Riemannian metric g . The elements of the complex $C^k(M, f)$ are generated by the critical points $q \in \text{Crit}(f)$ of the Morse function f , and grouped together by their index, $k = \text{ind}(q)$, the number of negative eigenvalues of the Hessian matrix of f at q . The differential of the complex requires the use of the metric and is given by the gradient flow, $-\nabla f$, from one critical point to another. We will assume throughout the paper that the gradient flow satisfy the Morse-Smale transversality condition, that is, the stable and unstable manifolds of any two critical points intersect transversely. The homology of the Morse complex is well-known to be isomorphic to the standard homology, and therefore, independent of the choice of the Morse function f and the metric g . As a corollary of this isomorphism, the Morse inequalities bound the Betti numbers of M in terms of the number of critical points of the Morse function.

We are interested here to consider Morse theory in the presence of a symplectic structure, that is, on a symplectic manifold (M^{2n}, ω) . On the cohomology side, besides the de Rham cohomology, Tsai, Tseng and Yau (TTY) [14, 16, 17] found novel symplectic cohomologies of differential forms. These cohomologies, which we will call TTY cohomologies and labelled by $F^p H(M, \omega)$, with $p = 0, 1, \dots, n - 1$, have interesting properties. For one, they can in general vary with the symplectic structure as seen in explicit examples of a six-dimensional nilmanifold [17] and of a three-ball product with a three-torus, $B^3 \times T^3$ [15]. These cohomologies have also been used to distinguish inequivalent symplectic structures on open 4-manifolds [6]. Of particular relevance here, when the symplectic structure is integral class $[\omega] \in H^2(M, \mathbb{Z})$, Tanaka and Tseng pointed out that the TTY cohomologies are isomorphic to the de Rham cohomologies of a higher dimensional sphere bundle over the symplectic manifold [12]. Specifically, denote by E_p the odd-dimensional sphere bundle $S^{2p+1} \rightarrow E_p \rightarrow M$ with Euler class ω^{p+1} , then $F^p H(M, \omega) \cong H_{dR}(E_p)$. Certainly, on this

sphere bundle, which is a smooth manifold, we can bound the dimensions of the de Rham cohomology $H_{dR}(E_p)$ and hence, $F^p H(M, \omega)$, by Morse or Morse-Bott inequalities.

To simplify the discussion, we will focus mostly in this paper on the $p = 0$ TTY cohomology, $PH(M)$, called *primitive* cohomology, and introduced in [17]. (The case of $p \geq 1$ can be straightforwardly generalized from the $p = 0$ case and will be described explicitly in the concluding section of this paper.) By [12], when ω is an integral class, $PH(M)$ are isomorphic to the de Rham cohomology of the prequantum circle bundle X , i.e. a circle bundle with Euler class given by ω . A bound on the dimension $PH(M) \cong H_{dR}(X)$ can be obtained by taking a Morse function f on M and pulling it back to the circle bundle, which makes $\pi^* f$ a Morse-Bott function. Denote by $b_k^\omega = \dim PH^k(M)$. The Morse-Bott inequalities for a circle bundle states the existence of a polynomial $Q(s)$ with positive coefficients such that

$$(1+s) \sum_{k=0}^{2n} m_k s^k = \sum_{k=0}^{2n+1} b_k^\omega s^k + (1+s)Q(s)$$

where m_k denotes the number of critical points with index equal to k . Specifically, this gives the strong inequalities

$$\sum_{i=0}^k (-1)^{k-i} b_i^\omega \leq \sum_{i=0}^k (-1)^{k-i} (m_i + m_{i-1}) = m_k \quad (1.1)$$

and the weak inequalities

$$b_k^\omega \leq m_k + m_{k-1}. \quad (1.2)$$

Though these Morse-Bott inequalities (1.1)-(1.2) assume ω is an integral class, they in fact hold true for any symplectic structure. We recall the algebraic relation for the primitive cohomologies in [14].

$$PH^k(M) \cong \text{coker} \left[\omega : H_{dR}^{k-2}(M) \rightarrow H_{dR}^k(M) \right] \oplus \ker \left[\omega : H_{dR}^{k-1}(M) \rightarrow H_{dR}^{k+1}(M) \right] \quad (1.3)$$

which immediately gives the weak inequalities of (1.2) just by bounding the dimensions of $H_{dR}^k(M)$ and $H_{dR}^{k-1}(M)$ by the number of Morse critical points, m_k and m_{k-1} , respectively. The strong inequalities can be similarly attained by applying the rank-nullity theorem. Hence, we find that the Morse-Bott inequalities only provide a rough estimate for b_k^ω .

Moreover, note that the b_k^ω 's on the left-hand-side of (1.1)-(1.2) generally depend on ω , while the m_k 's on the right hand side do not. These observations make us ask two questions:

- (1) Can we find a Morse-type complex that incorporate the symplectic structure ω explicitly and whose cohomology matches that of the TTY cohomology?
- (2) Can we bound $b_k^\omega = \dim PH^k(M, \omega)$ by inequalities that in general vary with ω ?

In this paper, we answer both questions in the affirmative. For the first, we introduce a symplectic Morse complex on (M^{2n}, ω) defined by a Morse-Smale pair (f, g) on M whose cohomology are isomorphic to the TTY primitive cohomology. Our symplectic Morse complex is motivated by the result of Tanaka-Tseng [12] which relates the cochain complex that underlies the TTY cohomologies with the cone complex of the wedge product map $\omega^{p+1} : \Omega^\bullet(M) \rightarrow \Omega^{\bullet+2p+2}(M)$ on the space of differential forms. Let us recall the definition of a cone on the de Rham complex with respect to ω^{p+1} . Again, for simplicity, we will focus on the case of $p = 0$.

Definition 1.1. Let (M^{2n}, ω) be a symplectic manifold. We define the **de Rham cone complex** of ω , $\text{Cone}(\omega) = (\Omega^\bullet(M) \oplus \theta \Omega^{\bullet-1}(M), d_C)$:

$$\dots \xrightarrow{d_C} \Omega^k(M) \oplus \theta \Omega^{k-1}(M) \xrightarrow{d_C} \Omega^{k+1}(M) \oplus \theta \Omega^k(M) \xrightarrow{d_C} \dots$$

where θ is a formal parameter of degree one and the differential d_C can be expressed in matrix form as

$$d_C = \begin{pmatrix} d & \omega \\ 0 & -d \end{pmatrix} \tag{1.4}$$

with d the standard exterior derivative and ω acting by wedge product.

Note that the d -closedness of ω together with the Leibniz product rule ensures that $d_C d_C = 0$. Also, if we formally define $d\theta = \omega$, then d_C is just the exterior derivative acting on $\Omega^\bullet(M) \oplus \theta \Omega^{\bullet-1}(M)$. Of interest, Tanaka-Tseng [12] proved the isomorphism of this cone cohomology with the TTY cohomology:

$$H(\text{Cone}(\omega)) \cong PH(M, \omega). \tag{1.5}$$

Motivated by the relationship between de Rham complex and the Morse cochain complex over \mathbb{R} , we define in the following a *cone* Morse complex with respect to ω also over \mathbb{R} .

Definition 1.2. Let (M^{2n}, ω) be a symplectic manifold equipped with a Riemannian metric g and a Morse function f satisfying the Morse-Smale transversality condition. Let $C^k(M, f)$ be the \mathbb{R} -module with generators the critical points of f with index k . We define the **cone Morse cochain complex** of ω , $\text{Cone}(c(\omega)) = (C^\bullet(M, f) \oplus C^{\bullet-1}(M, f), \partial_C)$:

$$\dots \xrightarrow{\partial_C} C^k(M, f) \oplus C^{k-1}(M, f) \xrightarrow{\partial_C} C^{k+1}(M, f) \oplus C^k(M, f) \xrightarrow{\partial_C} \dots$$

with

$$\partial_C = \begin{pmatrix} \partial & c(\omega) \\ 0 & -\partial \end{pmatrix}. \quad (1.6)$$

Here, ∂ is the standard Morse cochain differential defined by gradient flow, and $c(\omega) : C^k(M, f) \rightarrow C^{k+2}(M, f)$ acting on a critical point of index k is defined to be

$$c(\omega) q_k = \sum_{\text{ind}(r)=k+2} \left(\int_{\mathcal{M}(r_{k+2}, q_k)} \omega \right) r_{k+2} \quad (1.7)$$

where $\mathcal{M}(r_{k+2}, q_k)$ is the two-dimensional subspace of M consisting of all flow lines from the index $k+2$ critical point, r_{k+2} , to q_k .

Notice that the elements of the Morse cone complex, $\text{Cone}^k(c(\omega)) = C^k(M, f) \oplus C^{k-1}(M, f)$, can be generated by *pairs* of critical points of index k and $k-1$. The differential ∂_C consists of the standard Morse differential ∂ from gradient flow coupled with the $c(\omega)$ map which involves an integration of ω over the space of gradient flow lines. This type of maps has appeared in Austin-Braam [1] and Viterbo [18] to define a cup product on Morse cohomology. It satisfies the following commuting relation:

$$\partial(c(\omega)) = c(\omega)\partial. \quad (1.8)$$

(This relation follows from a more general Leibniz type formula. See [5, Appendix A].) With (1.8) and $\partial\partial = 0$, they together imply $\partial_C\partial_C = 0$. The cone Morse complex hence

gives the following cohomology for $k = 0, 1, \dots, 2n + 1$,

$$H^k(\text{Cone}(c(\omega))) = \frac{\ker \partial_C \cap \text{Cone}^k(c(\omega))}{\text{im } \partial_C \cap \text{Cone}^k(c(\omega))}.$$

For this cone Morse cohomology, we are able to prove the following theorem.

Theorem 1.3. *Let (M^{2n}, ω) be a closed symplectic manifold, The cohomology of symplectic Morse complex is isomorphic to the TTY cohomology, i.e. $H(\text{Cone}(c(\omega))) \cong PH(M)$.*

Theorem 1.3 importantly shows that the cohomology of the cone Morse complex is independent of the choice of both the Morse function and the Riemannian metric used to define $\text{Cone}(c(\omega))$. Moreover, the dependence on the symplectic structure is explicit in the differential ∂_C which involves the integration of ω over flow lines between pairs of critical points of the Morse function.

We shall prove Theorem 1.3 by means of the analytic method of Witten deformation, as developed in [2, 3, 20], to the $\text{Cone}(\omega)$ complex. Analogous to the Witten deformation of the de Rham complex [19], we deform the cone differential d_C and its adjoint d_C^* by the Morse function f parameterized by a real parameter $t \geq 0$:

$$d_{C,t} = e^{-tf} d_C e^{tf}, \quad d_{C,t}^* = e^{tf} d_C^* e^{-tf}.$$

This deformation is just a conjugation by e^{tf} , and hence, the cohomologies of the deformed cone complex $(\text{Cone}(\omega), d_{C,t})$ do not vary with t . However, the deformation has a significant effect on the harmonic forms, which provides an isomorphic description of the cohomologies. Specifically, the deformed cone Laplacian operator depends on t

$$\Delta_{C,t} = d_{C,t} d_{C,t}^* + d_{C,t}^* d_{C,t},$$

and in fact, its highest order t dependence is given by $t^2 \|df\|_x^2$. Hence, as $t \rightarrow \infty$, the harmonic forms must localize near $df = 0$, i.e. the critical points of f . This allows an identification, as $t \rightarrow \infty$, of each harmonic form of $\Delta_{C,t}$ with a critical point $p \in \text{Crit}(f)$. However, for finite t sufficiently large, some of the harmonic form in the limit of $t \rightarrow \infty$ becomes no longer harmonic. It is thus useful to consider the cone subspace $F_{C,t}^{[0,1]} \subset \text{Cone}(\omega)$ consisting of eigenforms of $\Delta_{C,t}$ with eigenvalue $\lambda \in [0, 1]$. Note that $(F_{C,t}^{[0,1]}, d_{C,t})$ is also a cochain complex since, $[\Delta_{C,t}, d_{C,t}] = 0$. Of interest, for t sufficiently large, the number of generators of $F_{C,t}^{[0,1]}$ matches exactly that of $\text{Cone}(c(\omega))$. This leads us to establish an isomorphism

between the cohomologies $H(\text{Cone}(c(\omega))) \cong H(F_{C,t}^{[0,1]}) \cong H(\text{Cone}(\omega)) \cong PH(M)$, proving Theorem 1.3. This successfully addresses the first question.

For the second question, we use the established quasi-isomorphism to obtain Morse-type inequalities for $H(\text{Cone}(\omega)) \cong PH(M)$. With m_k denoting the number of Morse critical points with index k , we are able to prove the following:

Theorem 1.4. *Let (M, ω, f, g) be a closed, symplectic manifold with the Morse function f and the Riemannian metric g satisfying the Morse-Smale transversality condition. Then there exists a polynomial $Q(s)$ with non-negative integer coefficients such that*

$$(1+s) \sum_{k=0}^{2n} m_k s^k - (s+s^2) \sum_{k=0}^{2n-2} v_k s^k = \sum_{k=0}^{2n+1} b_k^\omega s^k + (1+s)Q(s), \quad (1.9)$$

where $b_k^\omega = \dim H^k(\text{Cone}(\omega))$ and $v_k = \text{rank}(c(\omega) : C^k(M, f) \rightarrow C^{k+2}(M, f))$.

Alternatively, we have the following Morse-type inequalities:

(A) *Weak cone Morse inequalities*

$$b_k^\omega \leq m_k - v_{k-2} + m_{k-1} - v_{k-1}, \quad k = 0, \dots, 2n+1; \quad (1.10)$$

(B) *Strong cone Morse inequalities*

$$\sum_{i=0}^k (-1)^{k-i} b_i^\omega \leq m_k - v_{k-1}, \quad k = 0, \dots, 2n+1. \quad (1.11)$$

Furthermore, the above inequalities become equalities when the Morse function f is perfect, i.e. the Betti numbers $b_k = \dim H_{dR}^k(M) = m_k$ for all $k = 0, \dots, 2n$.

As was our goal for the second question, our cone Morse inequalities (4.1)-(4.2) can certainly vary with the symplectic structure ω . Specifically, the b_k^ω 's on the left-hand side and the v_k 's on the right-hand side both are defined with dependence on ω . This is in contrast to the m_k 's which are fixed by the choice of the Morse function f on M .

Beyond addressing our two main questions, let us point out that our analytic study of the symplectic cone Morse theory should be extendable to analyze symplectic manifolds with group actions. In particular, it would be interesting to study manifolds with hamiltonian group action and work out its equivariant or more general invariant cone cohomologies and also their corresponding cone Morse theory. We note of a recent work

[21] that used Witten deformation to study torus actions and their related invariant cohomologies. Moreover, the Witten deformation method has been successfully applied (see, for example [8, 13]) to prove Guillemin-Sternberg's conjecture [7] concerning the commutativity of symplectic reduction and geometric quantization. It is interesting to ask how symplectic reduction affect the TTY cohomologies $F^p H(M, \omega)$, or more directly its cone equivalent $H(\text{Cone}(\omega^{p+1}))$, and consider a different type of quantization of (M, ω) making use of the symplectic cone complex. This paper represents a first step in addressing these other interesting questions.

Finally, we mention that the de Rham cone complex and its Morse theory defined here with respect to the symplectic structure ω has a generalization that can be studied in a very general context. In a companion paper [5], we describe a general cone complex and its cone Morse theory on any oriented manifold M , with respect to any degree ℓ form $\psi \in \Omega^\ell(M)$ that is d -closed. It is a challenge to carry out the analytic Witten deformation method in this general setting, and hence, the discussion in [5] utilizes purely algebraic methods.

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2 Witten deformation method

We apply the Witten deformation method (see for example, Zhang's book [20]) to analyze the cone complex $\text{Cone}(\omega)$. In this section, we will introduce the deformed cone Laplacian and analyze its harmonic solutions and give a bound for the eigenvalues of non-harmonic eigenforms. We shall begin first with some preliminaries and also introduce our notations.

2.1 Cone Laplacian and its deformation

Let (M^{2n}, ω) be a closed symplectic manifold and let g be a compatible Riemannian metric. For $\eta_k, \eta'_k \in \Omega^k(M)$, we have the standard inner product,

$$\langle \eta_k, \eta'_k \rangle = \int_M \eta_k \wedge * \eta'_k, \quad (2.1)$$

where $*$: $\Omega^k(M) \rightarrow \Omega^{2n-k}(M)$ is the Hodge star operator. As introduced in the Introduction, we are interested in the cone forms with respect to the $\omega \wedge$ map:

$$\text{Cone}^k(\omega)(M) = \Omega^k(M) \oplus \theta \Omega^{k-1}(M) = \left\{ \eta_k + \theta \xi_{k-1} \mid \eta_k \in \Omega^k(M), \xi_{k-1} \in \Omega^{k-1}(M) \right\},$$

with $k = 0, 1, 2, \dots, 2n + 1$, and θ should be thought of as a formal one-form parameter, with the following two properties: (1) $d\theta = \omega$ and (2) $\theta \wedge \theta = 0$. The cone forms are essentially a pair of differential forms, and so the standard inner product (2.1) on M can be used to define a natural inner product on $\text{Cone}^*(\omega)(M)$,

$$\langle \eta_k + \theta \xi_{k-1}, \eta'_k + \theta \xi'_{k-1} \rangle_C = \langle \eta_k, \eta'_k \rangle + \langle \xi_{k-1}, \xi'_{k-1} \rangle. \quad (2.2)$$

This cone inner product can also be expressed in terms of a Hodge star-type operator. We define $*_C : \text{Cone}^k(\omega) \rightarrow \text{Cone}^{2n+1-k}(\omega)$ by

$$*_C (\eta_k + \theta \xi_{k-1}) = *_C \xi_{k-1} + \theta (-1)^k *_C \eta_k. \quad (2.3)$$

We can then write

$$\begin{aligned} \langle \eta_k + \theta \xi_{k-1}, \eta'_k + \theta \xi'_{k-1} \rangle_C &= \int_M \frac{\partial}{\partial \theta} ((\eta_k + \theta \xi_{k-1}) \wedge *_C (\eta'_k + \theta \xi'_{k-1})) \\ &= \int_M \eta_k \wedge *_C \eta'_k + \xi_{k-1} \wedge *_C \xi'_{k-1} \end{aligned}$$

where the derivative $(\partial/\partial\theta)$ satisfies $(\partial/\partial\theta)(\theta\eta_k) = \eta_k$ for any $\eta_k \in \Omega^k(M)$. (For ease of notation, when it is clear that we are considering the inner product for cone forms, we will simply write \langle , \rangle to denote the cone inner product of (2.2) and leave out the C subscript.)

Turning to the differential operators acting on $\text{Cone}^*(\omega)$, the differential, $d_C : \text{Cone}^k(\omega) \rightarrow \text{Cone}^{k+1}(\omega)$ is defined to be

$$d_C(\eta_k + \theta \xi_{k-1}) = d\eta_k + \omega \wedge \xi_{k-1} - \theta d\xi_{k-1},$$

which corresponds simply to the exterior derivative acting on $(\eta_k + \theta \xi_{k-1})$ and using $d\theta = \omega$. Clearly, $d_C d_C = 0$. And with respect to the cone inner product (2.2), the adjoint of d_C

has the form

$$\begin{aligned} d_C^*(\eta_k + \theta\xi_{k-1}) &= d^*\eta_k + \theta(\Lambda\eta_k - d^*\xi_{k-1}) \\ &= (-1)^k *_C d_C *_C (\eta_k + \theta\xi_{k-1}) \end{aligned} \quad (2.4)$$

where $\Lambda = \omega^*$ denotes the adjoint of ω with respect to the inner product (2.1) on M and has the expression $\Lambda = (-1)^k *_C \omega^*$ when acting on a k -form. For convenience, we will often express the cone form as a two-vector, $\sigma_k = \begin{pmatrix} \eta_k \\ \xi_{k-1} \end{pmatrix} \in \text{Cone}^k(\omega)$. In this notation, d_C and its adjoint d_C^* have the following matrix form

$$d_C = \begin{pmatrix} d & \omega \\ 0 & -d \end{pmatrix}, \quad d_C^* = \begin{pmatrix} d^* & 0 \\ \Lambda & -d^* \end{pmatrix}.$$

As in Witten [19], we deform the above differential operators by a Morse function f parameterized by a non-negative real number $t \in \mathbb{R}^+$. First, the deformed exterior derivative and its adjoint take the form

$$d_t = e^{-tf} d e^{tf} = d + t df, \quad d_t^* = e^{tf} d^* e^{-tf} = d^* + t \iota_{\nabla f}. \quad (2.5)$$

As for the cone differentials (d_C, d_C^*) , their deformation can be expressed simply in terms of (d_t, d_t^*) as follows,

$$d_{C,t} = e^{-tf} d_C e^{tf} = \begin{pmatrix} d_t & \omega \\ 0 & -d_t \end{pmatrix}, \quad d_{C,t}^* = e^{tf} d_C^* e^{-tf} = \begin{pmatrix} d_t^* & 0 \\ \Lambda & -d_t^* \end{pmatrix}. \quad (2.6)$$

Let us now turn our attention to the Laplacian operator associated to the de Rham cone complex:

$$\Delta_C = d_C d_C^* + d_C^* d_C = \begin{pmatrix} \Delta + \omega\Lambda & -d\Lambda + \Lambda d \\ -\omega d^* + d^* \omega & \Delta + \Lambda\omega \end{pmatrix}, \quad (2.7)$$

where $\Delta = dd^* + d^*d$ is the standard de Rham Laplacian. Under deformation by a Morse function f , it becomes

$$\Delta_{C,t} = d_{C,t} d_{C,t}^* + d_{C,t}^* d_{C,t} = \begin{pmatrix} \Delta_t + \omega\Lambda & -d_t^{\Lambda^*} \\ -d_t^\Lambda & \Delta_t + \Lambda\omega \end{pmatrix}, \quad (2.8)$$

where

$$\begin{aligned}\Delta_t &= d_t d_t^* + d_t^* d_t = \Delta + t(\mathcal{L}_{\nabla f} + \mathcal{L}_{\nabla f}^*) + t^2 \|df\|_x^2, \\ d_t^\Lambda &= d_t \Lambda - \Lambda d_t = d\Lambda - \Lambda d + t(df\Lambda - \Lambda df) = d^\Lambda + t(df\Lambda - \Lambda df), \\ d_t^{\Lambda^*} &= \omega d_t^* - d_t^* \omega = \omega d^* - d^* \omega + t(\omega \iota_{\nabla f} - \iota_{\nabla f} \omega) = d^{\Lambda^*} + t(\omega \iota_{\nabla f} - \iota_{\nabla f} \omega).\end{aligned}$$

In the above calculation, we used the notation $d^\Lambda := d\Lambda - \Lambda d$, and d^{Λ^*} is its adjoint. Furthermore, $\mathcal{L}_{\nabla f}$ is the Lie derivative with respect to the gradient of f , with $\mathcal{L}_{\nabla f}^*$ being its adjoint, and $\|df\|_x^2 = g^{ij} \partial_i f \partial_j f$ is the pointwise norm of df with respect to the Hodge metric on forms. We will call Δ_C the cone Laplacian and $\Delta_{C,t}$ the deformed cone Laplacian.

It turns out to be useful to also consider deforming the differential operators with respect to $(-f)$, which is also a Morse function. We will denote operators deformed by $(-f)$ with a $(-t)$ subscript. For instance,

$$d_{C,-t} := e^{-t(-f)} d_C e^{t(-f)} = \begin{pmatrix} d_{-t} & \omega \\ 0 & -d_{-t} \end{pmatrix}, \quad d_{C,-t}^* := e^{t(-f)} d_C^* e^{-t(-f)} = \begin{pmatrix} d_{-t}^* & 0 \\ \Lambda & -d_{-t}^* \end{pmatrix}.$$

and similarly, $\Delta_{C,-t} := d_{C,-t} d_{C,-t}^* + d_{C,-t}^* d_{C,-t}$. Noting from (2.3) that $*_C *_C = Id$ and from (2.4) that $d_C^* = (-1)^k *_C d_C *_C$ acting on $\sigma_k \in \text{Cone}^k(\omega)$, we find the following relations:

$$d_{C,-t} = (-1)^{k+1} *_C d_{C,t}^* *_C, \quad d_{C,-t}^* = (-1)^k *_C d_{C,t} *_C, \quad (2.9)$$

acting on $\sigma_k \in \text{Cone}^k(\omega)$. These immediately imply the following:

Lemma 2.1. *The cone form $\sigma = \begin{pmatrix} \eta_k \\ \xi_{k-1} \end{pmatrix} \in \text{Cone}^k(\omega)$ is a harmonic solution of $\Delta_{C,-t}$*

*for the Morse function $(-f)$, if and only if $*_C \sigma = \begin{pmatrix} *\xi_{k-1} \\ (-1)^k *\eta_k \end{pmatrix} \in \text{Cone}^{2n+1-k}(\omega)$ is a harmonic solution of $\Delta_{C,t}$ for the Morse function f . In particular, the harmonic conditions $d_{C,-t} \sigma = 0$ and $d_{C,-t}^* \sigma = 0$ hold if and only if $d_{C,t} (*_C \sigma) = 0$ and $d_{C,t}^* (*_C \sigma) = 0$.*

The relation also extends to all eigenforms of the deformed cone Laplacian.

Lemma 2.2. *The cone form $\sigma \in \text{Cone}^k(\omega)$ is an eigenform of $\Delta_{C,-t}$ if and only if $(*_C \sigma_k) \in \text{Cone}^{2n+1-k}(\omega)$ is an eigenform of $\Delta_{C,t}$.*

Proof. The statement follows from (2.9). Specifically, acting on $\sigma_k \in \text{Cone}^k(\omega)$, we have

$$\begin{aligned}\Delta_{C,-t} &= d_{C,-t}d_{C,-t}^* + d_{C,-t}^*d_{C,-t} \\ &= \left[(-1)^k *_C d_{C,t}^* *_C\right] \left[(-1)^k *_C d_{C,t} *_C\right] + \left[(-1)^{k-1} *_C d_{C,t} *_C\right] \left[(-1)^{k+1} *_C d_{C,t}^* *_C\right] \\ &= *_C \Delta_{C,t} *_C\end{aligned}$$

□

2.2 Local harmonic solutions of the deformed cone Laplacian

We are interested in studying the spectrum of the deformed cone Laplacian $\Delta_{C,t}$ when t is large. Following Witten's observation in [19], as $t \rightarrow \infty$, the order t^2 term of $\Delta_{C,t}$ in (2.8), $t^2 \|df\|_x^2$, dominates, and so eigenforms of $\Delta_{C,t}$ must localize around the critical points of the Morse function f . This localization greatly simplifies the study of the $\Delta_{C,t}$ spectrum and it turns out understanding the local eigenforms of $\Delta_{C,t}$ in a local chart around a critical point of f is sufficient for obtaining the Morse-type inequalities for $H(\text{Cone}(\omega))$. In this subsection, we will write down the local harmonic solutions for $\Delta_{C,t}$ for large t around critical points of f . The local non-harmonic eigenforms will be taken up in the next subsection.

For studying the eigenforms of $\Delta_{C,t}$ in a local neighborhood around $p \in \text{Crit}(f)$, it is useful to work in a local coordinate chart where (ω, g, f) all have standard canonical forms. For this, Stratmann [11] showed that in the neighborhood of any $p \in \text{Crit}(f)$, the pullback of f under a properly chosen symplectomorphism can be expressed in the following form,

$$f = f(p) + \sum_{\ell=1}^{n_f(p)} -x_\ell^2/2 + \sum_{\ell=n_f(p)+1}^{2n} x_\ell^2/2.$$

where $n_f(p)$ denotes the index at the critical point p . (We shall often use the notation $n_f(p)$ instead of $\text{ind}(p)$ to emphasize the dependence of the index on f .) In other words, it is possible to modify the Morse function via a series of local symplectomorphisms if needed so that around any $p \in \text{Crit}(f)$ we have what we shall call a compatible coordinate chart.

Definition 2.3. A local coordinate chart $\{x_j\}_{j=1,\dots,2n}$ around a critical point $p \in \text{Crit}(f)$ is called a **compatible coordinate chart** if the following properties are satisfied simultaneously:

- Darboux coordinates, i.e. $\omega = \sum_{i=1}^n dx_i \wedge dx_{i+n}$;
- Normal coordinates, i.e. $g_{ij}(x) = \delta_{ij} + O(|x|^2)$;
- Morse coordinates, $f = f(p) + \sum_{\ell=1}^{n_f(p)} -x_\ell^2/2 + \sum_{\ell=n_f(p)+1}^{2n} x_\ell^2/2$.

In the following, when discussing localized eigenform solutions of $\Delta_{C,t}$, we will by default work in such a compatible coordinate chart.

For local harmonic solutions $\sigma_k = \begin{pmatrix} \eta_k \\ \xi_{k-1} \end{pmatrix} \in \text{Cone}^k(\omega)$, the harmonic conditions $d_{C,t} \sigma_k = d_{C,t}^* \sigma_k = 0$, with $(d_{C,t}, d_{C,t}^*)$ given in (2.6), impose the following four conditions:

$$d_{C,t} \sigma_k = 0 : \quad (a) \ d_t \eta_k + \omega \wedge \xi_{k-1} = 0, \quad (b) \ d_t \xi_{k-1} = 0; \quad (2.10)$$

$$d_{C,t}^* \sigma_k = 0 : \quad (c) \ d_t^* \eta_k = 0, \quad (d) \ d_t^* \xi_{k-1} - \Lambda \eta_k = 0. \quad (2.11)$$

Notice that if both ω and Λ were set to zero, then the above four conditions become just $d_t \eta_k = d_t^* \eta_k = 0$, and $d_t \xi_{k-1} = d_t^* \xi_{k-1} = 0$, which are just the usual deformed harmonic conditions. Such solutions were described by Witten [19] which we recall here.

Lemma 2.4 (Witten [19]). *Around a critical point $p \in \text{Crit}(f)$ with index $n_f(p) = k$ described by a local compatible coordinate chart $\{x_i\}_{i=1, \dots, 2n}$, such that $g_{ij} = \delta_{ij}$ and*

$$f(x) = f(p) + \frac{1}{2} (-x_1^2 - \dots - x_k^2 + x_{k+1}^2 + \dots x_{2n}^2),$$

there is a one-dimensional k -form solution generated by

$$\zeta_k := e^{-t|x|^2/2} dx_1 \wedge \dots \wedge dx_k, \quad (2.12)$$

that satisfies the deformed harmonic conditions, i.e. $d_t \zeta_k = d_t^ \zeta_k = 0$.*

Hence, in the case where $\omega = 0$, we would have two types of localized harmonic solutions for $\sigma_k \in \text{Cone}^k(\omega)$, generated by $(\zeta_k, 0)$ at all $p \in \text{Crit}(f)$ with $n_f(p) = k$, and also, $(0, \zeta_{k-1})$ at all $q \in \text{Crit}(f)$ with $n_f(q) = k - 1$. It turns out that the existence of two types of harmonic generators for each cone degree k persists even when $\omega \neq 0$.

Proposition 2.5 (Local harmonic solutions of the deformed cone Laplacian). *For $\text{Cone}^k(\omega)$, there exist two types of local harmonic solutions of the deformed cone Laplacian $\Delta_{C,t}$: localized about critical points $p \in \text{Crit}(f)$ with index $n_f(p) = k$ and localized about critical points $q \in \text{Crit}(f)$ with index $n_f(q) = k - 1$. Each index k or $k - 1$ critical point has one generating harmonic solution.*

Proof. The deformed cone harmonic solutions must satisfy the four conditions in (2.10)-(2.11). We will describe the solutions first in the case for cone forms of degree $k \leq n$ and then for the case of degree $k \geq n + 1$.

Case (I): $0 \leq k \leq n$

When $k \leq n$, the harmonic solutions of $\Delta_{C,t}$ localized around critical points $p \in \text{Crit}(f)$ with index $n_f(p) = k$ are the standard type in (2.12). Expressed in compatible coordinates they are generated by

$$\begin{pmatrix} \eta_k \\ \xi_{k-1} \end{pmatrix} = \begin{pmatrix} \zeta_k \\ 0 \end{pmatrix} = \begin{pmatrix} e^{-t|x|^2/2} dx_1 \wedge \dots \wedge dx_k \\ 0 \end{pmatrix}. \quad (2.13)$$

Since $d_t \zeta_k = d_t^* \zeta_k = 0$, it remains only to check that $\Lambda \zeta_k = 0$. For this, we note that in a compatible coordinate chart, Λ takes the simple form

$$\Lambda = \sum_{i=1}^n \iota_{\partial_{x_{i+n}}} \iota_{\partial_{x_i}}. \quad (2.14)$$

Clearly then, we have $\Lambda \zeta_k = 0$ as long as $k \leq n$.

We describe now the localized harmonic solutions of $\Delta_{C,t}$ around critical points $q \in \text{Crit}(f)$ with index $n_f(p) = k - 1$. In the compatible coordinate chart, the Morse function takes the form

$$f(x) = f(q) + \frac{1}{2}(-x_1^2 - x_2^2 - \dots - x_{k-1}^2 + x_k^2 + \dots + x_{2n}^2), \quad (2.15)$$

and we also introduce a local one form

$$\tau = \sum_{i=1}^{k-1} -x_{i+n} dx_i + \frac{1}{2} \sum_{i=k}^n (x_i dx_{i+n} - x_{i+n} dx_i),$$

with the property $d\tau = \omega$. Then the following generates harmonic solutions around critical

points q with index $n_f(q) = k - 1$:

$$\begin{pmatrix} \eta_k \\ \xi_{k-1} \end{pmatrix} = \begin{pmatrix} -\tau \wedge \zeta_{k-1} \\ \zeta_{k-1} \end{pmatrix} = \begin{pmatrix} \frac{1}{2} e^{-t|x|^2/2} \sum_{i=k}^n (x_{i+n} dx_i - x_i dx_{i+n}) \wedge dx_1 \wedge \dots \wedge dx_{k-1} \\ e^{-t|x|^2/2} dx_1 \wedge \dots \wedge dx_{k-1} \end{pmatrix}. \quad (2.16)$$

Let us check conditions (a)-(d) in (2.10)-(2.11). For (a), it follows by direct computation that

$$d_t \eta_k = d_t(-\tau \wedge \zeta_{k-1}) = -d\tau \wedge \zeta_{k-1} + \tau \wedge d_t \zeta_{k-1} = -\omega \wedge \xi_{k-1}.$$

Condition (b) is trivially satisfied since $d_t \zeta_{k-1} = 0$. For (c), we use the expression $d_t^* = e^{tf} d^* e^{-tf}$ from (2.5) and compute also making use of (2.15):

$$\begin{aligned} d_t^* \eta_k &= e^{tf} d^* (e^{-tf} \eta_k) \\ &= \frac{1}{2} e^{tf} d^* \left(e^{-tf} e^{-t|x|^2/2} \sum_{i=k}^n (x_{i+n} dx_i - x_i dx_{i+n}) \wedge dx_1 \wedge \dots \wedge dx_{k-1} \right) \\ &= \frac{1}{2} e^{t(f-f(q))} d^* \left(\exp \left(-t \sum_{\ell=k}^{2n} x_\ell^2 \right) \sum_{i=k}^n (x_{i+n} dx_i - x_i dx_{i+n}) \wedge dx_1 \wedge \dots \wedge dx_{k-1} \right) \\ &= e^{t(f-f(q))} \exp \left(-t \sum_{\ell=k}^n x_\ell^2 \right) \left(t \sum_{i=k}^n (x_i x_{i+n} - x_{i+n} x_i) \right) dx_1 \wedge \dots \wedge dx_{k-1} = 0. \end{aligned}$$

And for condition (d), since $d_t^* \xi_{k-1} = d_t^* \zeta_{k-1} = 0$, we only need to check that $\Lambda \eta_k = 0$, which follows from using the expression of Λ in (2.14).

Case (II): $n + 1 \leq k \leq 2n + 1$

Using Lemma 2.1, the generators of the harmonic solutions for $k \geq n + 1$ can be straightforwardly obtained by applying $*_C$ to the generators found for $k \leq n$ in Case (I). Specifically, let $\tilde{\sigma}_j \in \text{Cone}^j(\omega)$ with $j \leq n$ be a harmonic solution of the deformed cone Laplacian with respect to a Morse function $(-f)$, i.e. $d_{C,-t} \tilde{\sigma}_j = d_{C,-t}^* \tilde{\sigma}_j = 0$. Then, by Lemma 2.1, $*_C \tilde{\sigma}_j \in \text{Cone}^k(\omega)$ with degree $k = 2n + 1 - j \geq n + 1$ is a harmonic solution with respect to f , i.e. $d_{C,t} (*_C \tilde{\sigma}_j) = d_{C,t}^* (*_C \tilde{\sigma}_j) = 0$. Thus, if we take $\tilde{\sigma}_j$ to be the two types of local harmonic generators (2.13) and (2.16) found for cone forms of degree $j \leq n$ with respect to a Morse function which we label by $-f$, then applying $*_C$ to them would

result in two types of harmonic generators for cone forms of degree $k = 2n + 1 - j$. And since $j = 0, 1, \dots, n$, this means $k = n + 1, \dots, 2n + 1$, as desired.

Let us comment about the minus sign difference in the Morse function for $\tilde{\sigma}_j$, which is associated with $(-f)$, versus $\sigma_k = *_C \tilde{\sigma}_j$, associated with f . Notice first that a critical point $p \in \text{Crit}(-f)$ of index $n_{-f}(p)$ would remain a critical point $p \in \text{Crit}(f)$ but with index $n_f(p) = 2n - n_{-f}(p)$. For in the compatible coordinate chart, we have

$$\begin{aligned} f &= f(p) + \frac{1}{2}(-x_1^2 - \dots - x_{n_f(p)}^2 + x_{n_f(p)+1}^2 + \dots + x_{2n}^2) \\ -f &= -f(p) + \frac{1}{2}(x_1^2 + \dots + x_{n_f(p)}^2 - x_{n_f(p)+1}^2 - \dots - x_{2n}^2) \\ &= -f(p) + \frac{1}{2}(-\tilde{x}_1^2 - \dots - \tilde{x}_{n_{-f}(p)}^2 + \tilde{x}_{n_{-f}(p)+1}^2 + \dots + \tilde{x}_{2n}^2) \end{aligned}$$

where in the last line, we have applied a change of coordinates $x_i \rightarrow \pm \tilde{x}_{2n+1-i}$ to return to compatible coordinate chart. (We note that the constant $f(p)$ is inconsequential in our discussion here since the deformed harmonic differentials $(d_{C,t}, d_{C,t}^*)$, given in (2.6), do not depend on it.) In particular, the two types of harmonic generators localized around index j critical points (2.13) and around index $j - 1$ critical points (2.16) of $(-f)$ will still be localized at the same critical points after applying $*_C$. But the index of the critical points would become $2n - j = k - 1$ and $2n - (j - 1) = k$, respectively, defined with respect to f .

We will write down explicitly the harmonic generators obtained from applying the $*_C$ map. Applying $*_C$ first to the generator in (2.13), we obtain the local harmonic solutions about any critical point q of index $n_f(q) = k - 1$:

$$\begin{pmatrix} \eta_k \\ \xi_{k-1} \end{pmatrix} = \begin{pmatrix} 0 \\ \zeta_{k-1} \end{pmatrix} = \begin{pmatrix} 0 \\ e^{-t|x|^2/2} dx_1 \wedge \dots \wedge dx_{k-1} \end{pmatrix}. \quad (2.17)$$

Applying $*_C$ to the generator given in (2.16), we obtain a second localized generator this time around any critical point p of index $n_f(p) = k$:

$$\begin{pmatrix} \eta_k \\ \xi_{k-1} \end{pmatrix} = \begin{pmatrix} \zeta_k \\ \iota_Z \zeta_k \end{pmatrix} = \begin{pmatrix} e^{-t|x|^2/2} dx_1 \wedge \dots \wedge dx_k \\ \frac{1}{2} e^{-t|x|^2/2} \frac{\omega^{r-1}}{(r-1)!} \wedge \sum_{i=1}^r (x_i dx_i + x_{n+i} dx_{n+i}) \wedge dx_{r+1} \wedge \dots \wedge dx_n \end{pmatrix}, \quad (2.18)$$

where $r = k - n$. In (2.18), we have noted that the second component can be expressed

in terms of an interior product by a vector which we have denoted by Z . To define this vector, note that by (2.3) and (2.16), the second component is given by

$$\begin{aligned}\xi_{k-1} &= (-1)^k * (-\tau \wedge e^{-t|x|^2/2} dx_{2n} \wedge \dots \wedge dx_{2n+1-k}) \\ &= \iota_Z \left(e^{-t|x|^2/2} dx_1 \wedge \dots \wedge dx_{k-1} \right)\end{aligned}$$

where $Z = -\tau'^{\sharp}$ is the musical isomorphism applied to the one form

$$\begin{aligned}\tau' &= \sum_{i=0}^{2n-k} x_{n-i} dx_{2n-i} - \frac{1}{2} \sum_{i=2n-k+1}^n x_{2n-i} dx_{n-i} - x_{n-i} dx_{2n-i} \\ \text{so } Z &= \sum_{i=0}^{2n-k} -x_{n-i} \frac{\partial}{\partial x_{2n-i}} + \frac{1}{2} \sum_{i=2n-k+1}^n x_{2n-i} \frac{\partial}{\partial x_{n-i}} - x_{n-i} \frac{\partial}{\partial x_{2n-i}}\end{aligned}$$

where τ' is τ under the symplectomorphism that maps $x_{2n+1-i} \rightarrow \pm x_i$ (to account for the sign change from $-f$ to f).

Finally, let us add that it can be checked directly that the above two harmonic generators for $k \geq n+1$ (2.17)-(2.18) satisfy the four deformed harmonic conditions (a)-(d) in (2.10)-(2.11) in a compatible coordinate chart. \square

Proposition 2.5 thus tells us that for each cone degree k , there exist at least two types of generators of local harmonic solutions of the deformed cone Laplacian. (In the next subsection, we will show that any other local eigenforms will have a non-zero eigenvalue and hence not harmonic.) Since the two types of harmonic generators are localized at critical points of index k and $k-1$, this implies each critical point $p \in \text{Crit}(f)$ of index $n_f(p)$ support two harmonic generators of cone degree $n_f(p)$ and $n_f(p)+1$. Below, we express this observation as a corollary.

Corollary 2.6. *Let $f : M \rightarrow \mathbb{R}$ be a Morse function on a symplectic manifold (M^{2n}, ω) . Then in the compatible coordinate chart around $p \in \text{Crit}(f)$ with index $n_f(p) = k$, where*

$$\omega = \sum_{i=1}^n dx_i \wedge dx_{i+n}, \quad g_{ij} = \delta_{ij}, \quad f = f(p) + \frac{1}{2} (-x_1^2 - \dots - x_k^2 + x_{k+1}^2 + \dots + x_{2n}^2),$$

we have the following two generators of harmonic solutions to the deformed cone Laplacian:

$$\begin{pmatrix} \zeta_k \\ \iota_Z \zeta_k \end{pmatrix} \in \text{Cone}^k(\omega), \quad \begin{pmatrix} -\tau \wedge \zeta_k \\ \zeta_k \end{pmatrix} \in \text{Cone}^{k+1}(\omega), \quad (2.19)$$

where $\zeta_k = e^{-t|x|^2/2} dx_1 \wedge \dots \wedge dx_k$, and

$$\tau = \sum_{i=1}^{n_f(p)} -x_{i+n} dx_i + \frac{1}{2} \sum_{i=n_f(p)+1}^n x_i dx_{i+n} - x_{i+n} dx_i, \quad 0 \leq n_f(p) \leq n-1 \quad (2.20)$$

$$\iota_Z = \sum_{i=0}^{2n-n_f(p)} -x_{n-i} \iota_{\partial_{2n-i}} + \frac{1}{2} \sum_{i=2n-n_f(p)+1}^n x_{2n-i} \iota_{\partial_{n-i}} - x_{n-1} \iota_{\partial_{2n-i}}, \quad n+1 \leq n_f(p) \leq 2n \quad (2.21)$$

and defining $\tau = 0$ when $n_f(p) \geq n$, and $\iota_Z = 0$ when $n_f(p) \leq n$.

2.3 Bounding eigenvalues of local eigenforms of the deformed cone Laplacian

Let $\begin{pmatrix} \eta_k \\ \xi_{k-1} \end{pmatrix}$ be any local eigenform of $\Delta_{C,t}$ with eigenvalue λ and not generated by the harmonic generators of the previous subsection. We wish to show that $\lambda > ct$ for some positive, non-zero constant c . Therefore, as t goes to infinity, λ grows to infinity at least of the order ct .

Theorem 2.7. *On a compatible local coordinate chart near a critical point p of a Morse function f , suppose $\begin{pmatrix} \eta_k \\ \xi_{k-1} \end{pmatrix}$ is orthogonal to the local harmonic solutions generated by (2.19). Then the following inequality holds,*

$$\left\langle \Delta_{C,t} \begin{pmatrix} \eta_k \\ \xi_{k-1} \end{pmatrix}, \begin{pmatrix} \eta_k \\ \xi_{k-1} \end{pmatrix} \right\rangle \geq ct \left\| \begin{pmatrix} \eta_k \\ \xi_{k-1} \end{pmatrix} \right\|^2 + O(\|\eta_k\|^2, \|\xi_{k-1}\|^2), \quad (2.22)$$

where c is a positive constant. In particular, the eigenvalue of any local non-harmonic eigenform is greater than $c't$ for some fixed constant $c' > 0$.

Proof. We start with the following computation,

$$\begin{aligned}
\left\langle \Delta_{C,t} \begin{pmatrix} \eta_k \\ \xi_{k-1} \end{pmatrix}, \begin{pmatrix} \eta_k \\ \xi_{k-1} \end{pmatrix} \right\rangle &= \left\langle \begin{pmatrix} \Delta_t + \omega\Lambda & -d_t^{\Lambda*} \\ -d_t^\Lambda & \Delta_t + \Lambda\omega \end{pmatrix} \begin{pmatrix} \eta_k \\ \xi_{k-1} \end{pmatrix}, \begin{pmatrix} \eta_k \\ \xi_{k-1} \end{pmatrix} \right\rangle \\
&= \left\langle \begin{pmatrix} \Delta_t \eta_k + \omega\Lambda \eta_k - d_t^{\Lambda*} \xi_{k-1} \\ -d_t^\Lambda \eta_k + \Delta_t \xi_{k-1} + \Lambda\omega \xi_{k-1} \end{pmatrix}, \begin{pmatrix} \eta_k \\ \xi_{k-1} \end{pmatrix} \right\rangle \\
&= \langle \Delta_t \eta_k, \eta_k \rangle + \langle \omega\Lambda \eta_k, \eta_k \rangle - \langle d_t^{\Lambda*} \xi_{k-1}, \eta_k \rangle + \langle \Delta_t \xi_{k-1}, \xi_{k-1} \rangle - \langle d_t^\Lambda \eta_k, \xi_{k-1} \rangle + \langle \Lambda\omega \xi_{k-1}, \xi_{k-1} \rangle \\
&= \langle \Delta_t \eta_k, \eta_k \rangle - 2\langle d_t^\Lambda \eta_k, \xi_{k-1} \rangle + \langle \Delta_t \xi_{k-1}, \xi_{k-1} \rangle + \|\Lambda \eta_k\|^2 + \|\omega \xi_{k-1}\|^2. \tag{2.23}
\end{aligned}$$

Note that Δ_t is self-adjoint and has a basis of orthogonal eigenforms of the form $\left(\prod_j H_{m_j}(\sqrt{t}x_j) \right) e^{-t|x|^2/2} dx_{I_k}$ where H_{m_j} is the m_j -th Hermite polynomial. The corresponding eigenvalues are $2t(\ell + \sum m_j) = 2t(\ell + m)$ where ℓ is the number of missing/different coordinates in $dx_{I_k} = dx_{i_1} \wedge \dots \wedge dx_{i_k}$ from harmonic form's $dx_1 \wedge \dots \wedge dx_{n_f}$ and $\sum m_j = m$ is the sum of the Hermite polynomial numbers. Imposing the L^2 -norm condition on η_k and ξ_{k-1} , we can express them as linear combinations of these eigenforms. For instance, we shall write

$$\eta_k = \sum_\alpha \eta^\alpha = \sum_\alpha a_\alpha \prod_j H_{m_j}(\sqrt{t}x_j) e^{-t|x|^2/2} dx_{I_k}, \tag{2.24}$$

where the index α denotes a certain combination of $\{I_k, \{m_j\}\}$, and similarly,

$$\xi_{k-1} = \sum_\beta \xi^\beta = \sum_\beta b_\beta \prod_j H_{m_j}(\sqrt{t}x_j) e^{-t|x|^2/2} dx_{I_{k-1}}, \tag{2.25}$$

with the degree of the form adjusted accordingly. Furthermore, we will refer to all η^α 's and ξ^β 's eigenforms as simple forms.

Below, we will prove the theorem assuming $k \leq n$. For $k > n$, we can use the observation of Lemma 2.2 that $\Delta_{C,-t} = *C \Delta_{C,t} *C$. Under $f \rightarrow -f$, we have the relation $n_f(p) = 2n - n_{-f}(p)$ and so a similar proof would follow.

As $\Delta_{C,t}$ is self-adjoint, we have that if $\begin{pmatrix} \eta_k \\ \xi_{k-1} \end{pmatrix}$ is orthogonal to both $\begin{pmatrix} \zeta_k \\ 0 \end{pmatrix}$ and

$$\begin{pmatrix} -\tau \wedge \zeta_{k-1} \\ \zeta_{k-1} \end{pmatrix},$$

$$\left\langle \begin{pmatrix} \eta_k \\ \xi_{k-1} \end{pmatrix}, \begin{pmatrix} \zeta_k \\ 0 \end{pmatrix} \right\rangle = 0 = \left\langle \begin{pmatrix} \eta_k \\ \xi_{k-1} \end{pmatrix}, \begin{pmatrix} -\tau \wedge \zeta_{k-1} \\ \zeta_{k-1} \end{pmatrix} \right\rangle$$

where ζ_k, ζ_{k-1} are solutions of Δ_t . We notice that η_k cannot have a nonzero component in the kernel of Δ_t , as then it would not be orthogonal to $\begin{pmatrix} \zeta_k \\ 0 \end{pmatrix}$. Note that ξ_{k-1} may have a component $b_\beta \zeta_{k-1}$ in the expansion of (2.25).

If ξ has a component of $\xi^\beta = b_\beta \zeta_{k-1}$, then we must have

$$\begin{aligned} 0 &= \left\langle \begin{pmatrix} \eta_k \\ \xi_{k-1} \end{pmatrix}, \begin{pmatrix} -\tau \wedge \zeta_{k-1} \\ \zeta_{k-1} \end{pmatrix} \right\rangle \\ &= \langle \eta_k, -\tau \wedge \zeta_{k-1} \rangle + b_\beta \langle \zeta_{k-1}, \zeta_{k-1} \rangle = \langle \eta_k, -\tau \wedge \zeta_{k-1} \rangle + b_\beta \|\zeta_{k-1}\|^2. \end{aligned}$$

Thus, there must be special components of η in the expansion (2.24), which we will label together by $\eta^{\alpha'}$, taking the following form

$$\eta^{\alpha'} = \sum_j (a_j x_{n+j} dx_j - a_{n+j} x_j dx_{n+j}) \wedge \zeta_{k-1} \quad \mathbf{and} \quad b_\beta \|\zeta_{k-1}\|^2 = -\langle \eta^{\alpha'}, -\tau \wedge \zeta_{k-1} \rangle. \quad (2.26)$$

Note that by normalization, we have

$$\|\zeta_{k-1}\|^2 = \int e^{-t|x|^2} dVol = \frac{\pi^n}{t^n},$$

$$-\langle \eta^{\alpha'}, -\tau \wedge \zeta_{k-1} \rangle = -\int \sum \frac{a_j}{2} x_j^2 e^{-t|x|^2} dVol = -\frac{\pi^n \sum a_j}{4t^{n+1}}. \quad (2.27)$$

Thus,

$$b_\beta = -\frac{\sum a_j}{4t}. \quad (2.28)$$

Next we calculate $d_t^\Lambda \eta^{\alpha'} = (d_t \Lambda - \Lambda d_t) \eta^{\alpha'}$. Note that a particular form

$$\chi_j = x_{n+j} e^{-t|x|^2/2} dx_j \wedge dx_1 \wedge \dots \wedge dx_{n_f}$$

has no components containing the wedge pair $dx_r \wedge dx_{r+n}$ for any $1 \leq r \leq n$, as $n_f \leq n$. Using the local form of Λ given in (2.14), we have $\Lambda\chi_j = 0$ and can compute $d_t^\Lambda \chi_j = (d_t \Lambda - \Lambda d_t)\chi_j$ as follows,

$$\begin{aligned}
(d_t \Lambda - \Lambda d_t)\chi_j &= -\Lambda d_t(x_{n+j} dx_j \zeta_{k-1}) \\
&= -\Lambda(d + tdf \wedge)(x_{n+j} dx_j \wedge \zeta_{k-1}) \\
&= -\Lambda\left(d(x_{n+j} dx_j) \wedge \zeta_{k-1} - x_{n+j} dx_j \wedge d\zeta_{k-1} - x_{n+j} dx_j (tdf \wedge \zeta_{k-1})\right) \\
&= -\Lambda\left(-dx_j \wedge dx_{n+j} \wedge \zeta_{k-1} - x_j dx_{n+j} \wedge d_t \zeta_{k-1}\right) \\
&= -\Lambda(-dx_j \wedge dx_{n+j} \wedge \zeta_{k-1} - x_j dx_{n+j} \wedge 0).
\end{aligned}$$

Hence, we have $d_t^\Lambda \chi_j = \Lambda(dx_j \wedge dx_{n+j} \zeta_{k-1})$. As ζ_{k-1} is primitive, this just gives $d_t^\Lambda \chi_j = \zeta_{k-1}$. A similar result for $\chi_{n+j} = -x_j dx_{n+j} \wedge \zeta_{k-1}$ also gives $d_t^\Lambda \chi_{n+j} = \zeta_{k-1}$. And adding all these terms together and using (2.28) gives

$$d_t^\Lambda \eta^{\alpha'} = \sum a_j \zeta_{k-1} = -4tb_\beta \zeta_{k-1}. \quad (2.29)$$

Plugging the above formula into (2.23), we get

$$\begin{aligned}
\left\langle \Delta_{C,t} \begin{pmatrix} \eta^{\alpha'} \\ b_\beta \zeta_{k-1} \end{pmatrix}, \begin{pmatrix} \eta^{\alpha'} \\ b_\beta \zeta_{k-1} \end{pmatrix} \right\rangle &= \lambda_{\eta'} \|\eta^{\alpha'}\|^2 - 2\langle d_t^\Lambda \eta^{\alpha'}, b_\beta \zeta_{k-1} \rangle + \|\Lambda \eta^{\alpha'}\|^2 + \|\omega b_\beta \zeta_{k-1}\|^2 \\
&= \lambda_{\eta'} \|\eta^{\alpha'}\|^2 - 2\langle -4tb_\beta \zeta_{k-1}, b_\beta \zeta_{k-1} \rangle + \|\Lambda \eta^{\alpha'}\|^2 + \|\omega b_\beta \zeta_{k-1}\|^2 \\
&= \lambda_{\eta'} \|\eta^{\alpha'}\|^2 + 8t \|b_\beta \zeta_{k-1}\|^2 + \|\Lambda \eta^{\alpha'}\|^2 + \|\omega b_\beta \zeta_{k-1}\|^2
\end{aligned}$$

Thus in this case, the inner product is bounded by ct with $c = 8$. We thus have the right estimate when ζ_{k-1} is harmonic and the specific $\eta^{\alpha'}$ that belongs to the span of $x_j dx_{n+j} \wedge \zeta_{k-1}$ and $x_{n+j} dx_j \wedge \zeta_{k-1}$ for $j = 1, \dots, n_f$.

Before we continue, we need to make some simple observations on the deformed Laplacian $\Delta_t^\Lambda = d_t^\Lambda d_t^{\Lambda*} + d_t^{\Lambda*} d_t^\Lambda$.

(i). It is useful to express $d_t^\Lambda = d_t \Lambda - \Lambda d_t = *_s d_t *_s$ in terms of the symplectic star operator $*_s$, which is defined analogous to the Hodge star operator but with respect to the symplectic structure instead of the Riemannian metric (for a reference, see [16, Sec. 2.1]).

And similarly, $d_t^{\Lambda*} = \omega d_t^* - d_t^* \omega = *_s d_t^* *_s$. Therefore, using that $*_s^2 = \mathbb{I}$, we have

$$\begin{aligned}\Delta_t^\Lambda &= d_t^\Lambda d_t^{\Lambda*} + d_t^{\Lambda*} d_t^\Lambda \\ &= *_s d_t *_s *_s d_t^* *_s + *_s d_t^* *_s *_s d_t *_s \\ &= *_s (d_t d_t^* + d_t^* d_t) *_s \\ &= *_s \Delta_t *_s\end{aligned}$$

Therefore, since $*_s = *_s^{-1}$, $\Delta_t^\Lambda = *_s \Delta_t *_s^{-1}$ has the eigenforms $*_s \prod_j H_{m_j}(\sqrt{t}x_j) e^{-t|x|^2/2} dx_{I_k} = \prod_j H_{m_j}(\sqrt{t}x_j) *_s dx_{I_k}$ with eigenvalues $2t(\ell_\Lambda + \sum m_j)$, where ℓ^Λ is the number of forms in $*_s dx_{I_k}$ missing/different from $dx_1 \wedge \dots \wedge dx_{n_f}$.

(ii). Note that

$$\begin{aligned}\|d_t^\Lambda \eta_k\|^2 &= \langle d_t^\Lambda \eta_k, d_t^\Lambda \eta_k \rangle \\ &\leq \langle d_t^\Lambda \eta_k, d_t^\Lambda \eta_k \rangle + \langle d_t^{\Lambda*} \eta_k, d_t^{\Lambda*} \eta_k \rangle \\ &= \langle d_t^{\Lambda*} d_t^\Lambda \eta_k, \eta_k \rangle + \langle d_t^\Lambda d_t^{\Lambda*} \eta_k, \eta_k \rangle \\ &= \langle \Delta_t^\Lambda \eta_k, \eta_k \rangle.\end{aligned}$$

A similar argument shows $\|d_t^{\Lambda*} \xi_{k-1}\|^2 \leq \langle \Delta_t^\Lambda \xi_{k-1}, \xi_{k-1} \rangle$.

With these inequalities, we now proceed to show that

$$\left\langle \Delta_{C,t} \begin{pmatrix} \eta_k \\ \xi_{k-1} \end{pmatrix}, \begin{pmatrix} \eta_k \\ \xi_{k-1} \end{pmatrix} \right\rangle \geq ct \left\| \begin{pmatrix} \eta_k \\ \xi_{k-1} \end{pmatrix} \right\|^2 + O(\|\eta_k\|^2, \|\xi_{k-1}\|^2).$$

Writing the forms in terms of linear combinations of eigenforms, $\eta_k = \sum_\alpha \eta^\alpha$ and $\xi_{k-1} = \sum_\beta \xi^\beta$, as in (2.24)-(2.25), the expression for the inner product in (2.23) becomes

$$\begin{aligned}&\left\langle \Delta_{C,t} \begin{pmatrix} \sum_\alpha \eta^\alpha \\ \sum_\beta \xi^\beta \end{pmatrix}, \begin{pmatrix} \sum_\alpha \eta^\alpha \\ \sum_\beta \xi^\beta \end{pmatrix} \right\rangle \\ &= \sum_\alpha \lambda_{\eta^\alpha} \langle \eta^\alpha, \eta^\alpha \rangle - \sum_{\alpha, \beta} 2 \langle d_t^\Lambda \eta^\alpha, \xi^\beta \rangle + \sum_\beta \lambda_{\xi^\beta} \langle \xi^\beta, \xi^\beta \rangle + \sum_{\alpha, \beta} O(\|\eta^\alpha\|^2, \|\xi^\beta\|^2)\end{aligned}\tag{2.30}$$

where $\|\Lambda\eta^\alpha\|^2$ and $\|\omega\xi^\beta\|^2$ are bounded by (constant multiple of) the norm of $\|\eta^\alpha\|^2$ and $\|\xi^\beta\|^2$. To estimate the $\langle d_t^\Lambda\eta^\alpha, \xi^\beta \rangle$ term, we start by applying the Cauchy-Schwarz inequality and arithmetic-geometric mean inequality to the pair $(\frac{1}{C}\|d_t^\Lambda\eta^\alpha\|^2, C\|\xi^\beta\|^2)$ for any positive constant C . Thus

$$\langle d_t^\Lambda\eta^\alpha, \xi^\beta \rangle \leq |\langle d_t^\Lambda\eta^\alpha, \xi^\beta \rangle| \leq \|d_t^\Lambda\eta^\alpha\| \cdot \|\xi^\beta\| \leq \frac{1}{2} \left(\frac{1}{C} \|d_t^\Lambda\eta^\alpha\|^2 + C \|\xi^\beta\|^2 \right). \quad (2.31)$$

Next, note that $\eta^\alpha = a_\alpha \prod_r H_{m_r}(\sqrt{t}x_r) e^{-t|x|^2/2} dx_{I_k}$ and so we can compute $d_t\eta^\alpha$ using the two standard properties of the Hermite polynomials:

$$\begin{aligned} \frac{\partial}{\partial x_r} H_{m_r}(\sqrt{t}x_r) &= \sqrt{t}(2m_r)H_{m_r-1}(\sqrt{t}x_r), \\ \sqrt{t}x_r H_{m_r}(\sqrt{t}x_r) &= \frac{1}{2}H_{m_r+1}(\sqrt{t}x_r) + m_r H_{m_r-1}(\sqrt{t}x_r). \end{aligned}$$

Writing $f = f(p) + \sum_j \nu_j \frac{x_j^2}{2}$, where $\nu_j = \pm 1$, and so $df = \sum_j \nu_j x_j dx_j$, we obtain

$$\begin{aligned} d_t\eta^\alpha &= (d + tdf)a_\alpha \prod_r H_{m_r}(\sqrt{t}x_r) e^{-t|x|^2/2} \wedge dx_{I_k} \\ &= a_\alpha \left[d \left(\prod_r H_{m_r}(\sqrt{t}x_r) e^{-t|x|^2/2} \right) + \sum_j \nu_j t x_j dx_j \prod_r H_{m_r}(\sqrt{t}x_r) e^{-t|x|^2/2} \right] \wedge dx_{I_k} \\ &= a_\alpha \left[\sum_s \left(\prod_{r \neq s} H_{m_r}(\sqrt{t}x_{m_r}) \left(\frac{\partial}{\partial x_s} H_{m_s}(\sqrt{t}x_s) \right) e^{-t|x|^2/2} \right) dx_s \right. \\ &\quad \left. + \sum_j (\nu_j - 1) t x_j H_{m_j}(\sqrt{t}x_j) e^{-t|x|^2/2} \prod_{r \neq j} H_{m_r}(\sqrt{t}x_r) dx_j \right] \wedge dx_{I_k} \\ &= a_\alpha \left[\sum_s \left(\prod_{r \neq s} H_{m_r}(\sqrt{t}x_{m_r}) \left(2\sqrt{t}m_s H_{m_s-1}(\sqrt{t}x_s) e^{-t|x|^2/2} \right) \right) dx_s \right. \\ &\quad \left. + \sum_j \sqrt{t}(\nu_j - 1) \left(\frac{1}{2} H_{m_j+1}(\sqrt{t}x_j) + m_j H_{m_j-1}(\sqrt{t}x_j) \right) \prod_{r \neq j} H_{m_r}(\sqrt{t}x_r) dx_j \right] \wedge dx_{I_k}. \end{aligned}$$

The expression for $d_t\eta^\alpha$ above contains three distinct terms with $H_{m_r \neq s} H_{m_s-1}$ associated with the basic form $dx_s \wedge dx_{I_k}$, and $H_{m_r \neq j} H_{m_j+1}$ and $H_{m_r \neq j} H_{m_j-1}$ both associated with

$dx_j \wedge dx_{I_k}$. These terms are summed over the $2n$ possible r , so we have at most $6n$ terms in $d_t \eta^\alpha$ when η^α is basic. Next, note that $\Lambda = \iota_{\partial_{r+n}} \iota_{\partial_r}$ can have at most n terms when applied to a basic η^α . Thus, $d_t^\Lambda \eta^\alpha = d_t \Lambda \eta^\alpha - \Lambda d_t \eta^\alpha$ can have $6n^2$ basic terms for $d_t \Lambda$ and $6n^2$ terms for Λd_t so we get at most $12n^2$ terms (some of these may cancel out). Thus, for each of the basic η^α , the term $\langle d_t^\Lambda \eta^\alpha, \xi^\beta \rangle$ is nonzero for at most $12n^2$ of the ξ^β . Let us label this set of at most $12n^2$ terms by

$$S_{\eta^\alpha} = \{\xi^\beta : \langle d_t^\Lambda \eta^\alpha, \xi^\beta \rangle \neq 0\}. \quad (2.32)$$

To bound our inequality, we look at a particular α in (2.30),

$$\lambda_{\eta^\alpha} \langle \eta^\alpha, \eta^\alpha \rangle - \sum_{\beta} 2 \langle d_t^\Lambda \eta^\alpha, \xi^\beta \rangle + \sum_{\beta} \lambda_{\xi^\beta} \langle \xi^\beta, \xi^\beta \rangle + O(\|\eta^\alpha\|^2, \|\xi^\beta\|^2).$$

Ignoring the non t term $O(\|\eta^\alpha\|^2, \|\xi^\beta\|^2)$, we need to examine the cross terms $-2 \langle d_t \eta^\alpha, \xi^\beta \rangle$. For a particular η^α , and all the ξ^β that have a nonzero contribution (the at most $12n^2$ we found above). Thus, if we look at the sum of those terms and use our (2.31), we have

$$\begin{aligned} \lambda_{\eta^\alpha} \|\eta^\alpha\|^2 - \sum_{\xi^\beta \in S_{\eta^\alpha}} 2 \langle d_t^\Lambda \eta^\alpha, \xi^\beta \rangle &\geq \lambda_{\eta^\alpha} \|\eta^\alpha\|^2 - \sum_{\xi^\beta \in S_{\eta^\alpha}} 2 |\langle d_t^\Lambda \eta^\alpha, \xi^\beta \rangle| \\ &\geq 2t(\ell^{\eta^\alpha} + m^{\eta^\alpha}) \|\eta^\alpha\|^2 - \left(\sum_{\xi^\beta \in S_{\eta^\alpha}} \frac{1}{C} \|\Delta_{C,t}^\Lambda \eta^\alpha\|^2 + C \|\xi^\beta\|^2 \right) \\ &\geq 2t(\ell^{\eta^\alpha} + m^{\eta^\alpha}) \|\eta^\alpha\|^2 - \sum_{\xi_\alpha \in S_{\eta^\alpha}} \frac{2t}{C} (\ell_\Lambda^{\eta^\alpha} + m^{\eta^\alpha}) \|\eta^\alpha\|^2 - C \|\xi^{\beta'}\|^2 \\ &= 2t \left(\ell^{\eta^\alpha} - \sum_{\xi_\alpha \in S_{\eta^\alpha}} \frac{\ell_\Lambda^{\eta^\alpha}}{C} + \left(1 - \sum_{\xi_\alpha \in S_{\eta^\alpha}} \frac{1}{C} \right) m^{\eta^\alpha} \right) \|\eta^\alpha\|^2 - \sum_{\xi^\beta \in S_{\eta^\alpha}} C \|\xi^\beta\|^2. \end{aligned}$$

Now, recall that η_k is not harmonic, so either $\sum m^{\eta^\alpha} \geq 1$ or $\ell^{\eta^\alpha} \geq 1$. Also, $\ell_\Lambda^{\eta^\alpha} \leq 2n$, as we can have at most $2n$ terms different from $dx_1 \wedge \dots \wedge dx_{n_f}$. Then we can bound the terms above with a positive constant by setting $12^2 n^3 = C > 6n \cdot (12n^2)$, and showing the following is positive

$$t \left(\ell^{\eta^\alpha} - \sum_{\xi_\alpha \in S_{\eta^\alpha}} \frac{\ell_\Lambda^{\eta^\alpha}}{C} + \left(1 - \sum_{\xi_\alpha \in S_{\eta^\alpha}} \frac{1}{C} \right) m^{\eta^\alpha} \right) \|\eta^\alpha\|^2.$$

We note that this is summed over S_{η^α} , which we know has at most $12n^2$ terms, and changing ξ^β does not change this portion of the sum, so

$$\begin{aligned}
& 2t \left(\ell^{\eta^\alpha} - \sum_{\xi_\alpha \in S_{\eta^\alpha}} \frac{\ell_\Lambda^{\eta^\alpha}}{C} + \left(1 - \sum_{\xi_\alpha \in S_{\eta^\alpha}} \frac{1}{C} \right) m^{\eta^\alpha} \right) \|\eta^\alpha\|^2 \\
& \geq 2t \left(\ell^{\eta^\alpha} - \frac{12n^2 \ell_\Lambda^{\eta^\alpha}}{C} + \left(1 - \frac{12n^2}{C} \right) m^{\eta^\alpha} \right) \|\eta^\alpha\|^2 \\
& = 2t \left(\ell^{\eta^\alpha} - \left(\frac{12n^2 \ell_\Lambda^{\eta^\alpha}}{12^2 n^3} \right) + \left(1 - \frac{12n^2}{12^2 n^3} \right) m^{\eta^\alpha} \right) \|\eta^\alpha\|^2 \\
& = 2t \left(\ell^{\eta^\alpha} - \frac{\ell_\Lambda^{\eta^\alpha}}{6n} + \left(1 - \left(\frac{1}{6n} \right) \right) m^{\eta^\alpha} \right) \|\eta^\alpha\|^2.
\end{aligned}$$

To investigate the above term, we split into two cases.

Case 1: $\ell^{\eta^\alpha} \geq 1$. Thus

$$\begin{aligned}
2t \left(\ell^{\eta^\alpha} - \frac{\ell_\Lambda^{\eta^\alpha}}{6n} + \left(1 - \left(\frac{1}{6n} \right) \right) m^{\eta^\alpha} \right) \|\eta^\alpha\|^2 & \geq 2t \left(\ell^{\eta^\alpha} - \frac{\ell_\Lambda^{\eta^\alpha}}{6n} + \left(1 - \left(\frac{1}{6n} \right) \right) 0 \right) \|\eta^\alpha\|^2 \\
& \geq 2t \left(1 - \frac{2n}{6n} \right) \|\eta^\alpha\|^2 \geq t \|\eta^\alpha\|^2.
\end{aligned}$$

Case 2: Suppose $\ell^{\eta^\alpha} \geq 0$ and $m^{\eta^\alpha} \geq 1$. Then we have

$$\begin{aligned}
2t \left(\ell^{\eta^\alpha} - \frac{\ell_\Lambda^{\eta^\alpha}}{6n} + \left(1 - \left(\frac{1}{6n} \right) \right) m^{\eta^\alpha} \right) \|\eta^\alpha\|^2 & \geq 2t \left(0 - \frac{2n}{6n} + \left(1 - \frac{1}{6} \right) (1) \right) \|\eta^\alpha\|^2 \\
& \geq 2t \left(\frac{-1}{3} + \frac{5}{6} \right) \|\eta^\alpha\|^2 \geq t \|\eta^\alpha\|^2.
\end{aligned}$$

Thus, we can conclude that c^{η^α} is positively bounded in all cases. So we have

$$\lambda_{\eta^\alpha} \|\eta^\alpha\|^2 - \sum_{\xi^\beta \in S_{\eta^\alpha}} 2 \langle d_t^\Lambda \eta^\alpha, \xi^\beta \rangle \geq t \|\eta^\alpha\|^2 - 12^2 n^3 |S_{\eta^\alpha}| \|\xi^\beta\| \geq t \|\eta^\alpha\|^2 - 12^2 n^3 (12n^2) \|\xi^\beta\| \tag{2.33}$$

where S_{η^α} has at most $12n^2$ elements. Now looking back at our original expression (2.30)

$$\lambda_{\eta^\alpha} \|\eta^\alpha\|^2 - 2 \langle d_t^\Lambda \eta^\alpha, \xi^\beta \rangle + \lambda_{\xi^\beta} \|\xi^\beta\|^2 + O(\|\eta^\alpha\|, \|\xi^\beta\|),$$

if we replace $2\langle d_t^\Lambda \eta^\alpha, \xi^\beta \rangle \leq 2|\langle d_t^\Lambda \eta^\alpha, \xi^\beta \rangle|$ and use our inequality above (which puts the cross terms with a particular η^α), we then obtain

$$\begin{aligned}
& \lambda_{\eta^\alpha} \|\eta^\alpha\|^2 - 2\langle d_t^\Lambda \eta^\alpha, \xi^\beta \rangle + \lambda_{\xi^\beta} \|\xi^\beta\|^2 + O(\|\eta^\alpha\|, \|\xi^\beta\|) \\
& \geq \lambda_{\eta^\alpha} \|\eta^\alpha\|^2 - 2|\langle d_t^\Lambda \eta^\alpha, \xi^\beta \rangle| + \lambda_{\xi^\beta} \|\xi^\beta\|^2 + O(\|\eta^\alpha\|, \|\xi^\beta\|) \\
& \geq t\|\eta^\alpha\|^2 - 12^2 n^3 \|\xi^\beta\|^2 + \lambda_{\xi^\beta} \|\xi^\beta\| + O(\|\eta^\alpha\|, \|\xi^\beta\|) \\
& \geq t\|\eta^\alpha\|^2 + \lambda_{\xi^\beta} \|\xi^\beta\| + O(\|\eta^\alpha\|, \|\xi^\beta\|),
\end{aligned}$$

where we have used the property that $12^2 n^3 \|\xi^\beta\|^2$ is independent of t .

Next let $\eta = \eta^{\alpha'} + \sum_{\alpha} \eta^\alpha$ and $\xi = \xi^{\beta'} + \sum_{\beta} \xi^\beta$, where $\xi^{\beta'} = b_\beta \zeta_{k-1}$, ξ^β are the non-harmonic components, and $\eta^{\alpha'} = \sum (a_j x_{n+j} dx_j - a_{n+j} x_j dx_{n+j}) \wedge \zeta_{k-1}$ is the component of η we know exists to be orthogonal to our two solutions by (2.26), and η^α are the components orthogonal to $\eta^{\alpha'}$. Note we showed that $\langle d_t^\Lambda \eta^{\alpha'}, \xi_\beta \rangle = 0$ for the $\xi^\beta \neq \xi^{\beta'}$ as these are orthogonal to $d_t^\Lambda \eta^{\alpha'} = b_\beta \zeta_{k-1}$. Using (2.29), (2.33), and (2.30) to both η_k and ξ_{k-1} we have

$$\begin{aligned}
\lambda_{C,t} \left\| \begin{array}{c} \eta_k \\ \xi_{k-1} \end{array} \right\|^2 &= \left\langle \Delta_{C,t} \begin{pmatrix} \eta^{\alpha'} + \sum_{\alpha} \eta^\alpha \\ \xi^{\beta'} + \sum_{\beta} \xi^\beta \end{pmatrix}, \begin{pmatrix} \eta^{\alpha'} + \sum_{\alpha} \eta^\alpha \\ \eta^{\beta'} + \sum_{\beta} \xi^\beta \end{pmatrix} \right\rangle \\
&\geq \lambda_{\eta^{\alpha'}} \|\eta^{\alpha'}\|^2 - 2\langle d_t^\Lambda \eta^{\alpha'}, \xi^{\beta'} \rangle + 0\|\xi^{\beta'}\|^2 + \sum_{\alpha} \lambda_{\eta^\alpha} \|\eta^\alpha\|^2 - \sum_{\alpha, \beta, \beta'} 2\langle d_t^\Lambda \eta^\alpha, \xi^\beta \rangle + \\
&\quad \sum_{\beta} \lambda_{\xi^\beta} \|\xi^\beta\|^2 + O(\|\eta^\alpha\|, \|\xi^\beta\|, \|\eta^{\alpha'}\|, \|\xi^{\beta'}\|) \\
&\geq \lambda_{\eta^{\alpha'}} \|\eta^{\alpha'}\|^2 + 8t\|\xi^{\beta'}\|^2 + \sum_{\alpha} \lambda_{\eta^\alpha} \|\eta^\alpha\|^2 - 2 \sum_{\xi^\beta \in S_{\eta^\alpha}} |\langle d_t^\Lambda \eta^{\alpha'}, \xi^\beta \rangle| + \\
&\quad \sum_{\beta} \lambda_{\xi^\beta} \|\xi^\beta\|^2 + O(\|\eta^\alpha\|, \|\xi^\beta\|, \|\eta^{\alpha'}\|, \|\xi^{\beta'}\|) \\
&\geq t\|\eta^{\alpha'}\|^2 + \sum_{\alpha} t\|\eta^\alpha\|^2 + t\|\xi^{\beta'}\|^2 - \sum_{\beta} 12^2 n^3 (12n^2) \|\xi^\beta\|^2 + \\
&\quad \sum_{\beta} t\|\xi^\beta\|^2 + O(\|\eta^\alpha\|, \|\xi^\beta\|, \|\eta^{\alpha'}\|, \|\xi^{\beta'}\|) \\
&\geq t \left\| \begin{array}{c} \eta^\alpha \\ \xi^\beta \end{array} \right\|^2 - O(\|\eta^\alpha\|, \|\eta^{\alpha'}\|, \|\xi^\beta\|, \|\xi^{\beta'}\|)
\end{aligned}$$

where we know each λ_{ξ_β} is bounded by $2t$. Thus each term in our (possibly infinite) sum has a scale of t while it is added/subtracted by elements of order $O(\|\eta^\alpha\|, \|\eta^{\alpha'}\|, \|\xi^\beta\|, \|\xi^{\beta'}\|)$ with no factor of t . So by driving t large enough it will dominate this inequality (as has no factor of t) and as eigenforms are non-zero $\lambda > ct$ for sufficiently large t .

As another result of our inequality above, we now show that the kernel of $\Delta_{C,t}$ is generated by only our two harmonic solutions. For if it were a third independent solution, we could project it to the orthogonal complement of our two solutions and get a nonzero harmonic solution in the orthogonal complement of $\begin{pmatrix} \zeta_k \\ 0 \end{pmatrix}, \begin{pmatrix} -\tau \wedge \zeta_{k-1} \\ \zeta_{k-1} \end{pmatrix}$, but then our inequality above would apply to our new solution, but choosing t large enough would give a positive eigenvalue, contradicting that our projection was harmonic. Thus the kernel of $\Delta_{C,t}$ is two dimensional and is generated by $\begin{pmatrix} \zeta_k \\ 0 \end{pmatrix}, \begin{pmatrix} -\tau \wedge \zeta_{k-1} \\ \zeta_{k-1} \end{pmatrix}$. \square

2.4 Local harmonic solutions approximating eigenforms of the deformed cone Laplacian

In this subsection, we will show that the local harmonic solutions found in Section 2.2 can approximate global eigenforms of $\Delta_{C,t}$ when t is large. In particular, as we make clear in Theorem 2.13, the local harmonic solutions represent all the low-lying eigenforms (i.e. those with small eigenvalues) when t is sufficiently large. The discussion here will parallel that for the Witten-deformed de Rham complex as described in [20, Section 5.6].

Without loss of generality, we assume that each V_p where $p \in \text{Crit}(f)$ is an open ball of radius $4a$, and assume $t > 0$. Let γ_p be a smooth bump function such that $\gamma_p(z) = 1$ for $|z| \leq a$ and $\gamma_p(z) = 0$ for $|z| \geq 2a$. Define

$$\Phi_{p,1}(t) = \sqrt{\int_{V_p} \gamma_p^2 \left\| \begin{pmatrix} \zeta_k \\ \iota_Z \zeta_k \end{pmatrix} \right\|_x^2 dVol}, \quad \Phi_{p,2}(t) = \sqrt{\int_{V_p} \gamma_p^2 \left\| \begin{pmatrix} -\tau \wedge \zeta_k \\ \zeta_k \end{pmatrix} \right\|_x^2 dVol},$$

where $\left\| \begin{pmatrix} \eta_k \\ \xi_{k-1} \end{pmatrix} \right\|_x$ is the norm given by $\left\langle \begin{pmatrix} \eta_k \\ \xi_{k-1} \end{pmatrix}, \begin{pmatrix} \eta'_k \\ \xi'_{k-1} \end{pmatrix} \right\rangle_x dVol = \eta_k \wedge * \eta'_k + \xi_{k-1} \wedge * \xi'_{k-1}$, the pointwise form inner product induced by $*$. We further define

$$\rho_{p,1}(t) = \frac{\gamma_p}{\Phi_{p,1}(t)} \begin{pmatrix} \zeta_k \\ \iota_Z \zeta_k \end{pmatrix}, \quad \rho_{p,2}(t) = \frac{\gamma_p}{\Phi_{p,2}(t)} \begin{pmatrix} -\tau \wedge \zeta_k \\ \zeta_k \end{pmatrix}, \quad (2.34)$$

which are unit norm with compact support contained in V_p . Note that $\rho_{p,1}(t), \rho_{p,2}(t)$ encompass both generators of local harmonic solutions around $p \in \text{Crit}(f)$ described in Corollary 2.6.

Now let $\mathbf{H}^1(\text{Cone}(\omega))$ be the first Sobolev space with respect to a Sobolev norm on $\text{Cone}(\omega)$. Let $E_{C,t}$ denote the direct sum of the vector spaces generated by the $\rho_{p,1}(t), \rho_{p,2}(t)$, and let $E_{C,t}^\perp$ be the orthogonal complement to $E_{C,t}$ in $\mathbf{H}^1(\text{Cone}(\omega))$, so that $\mathbf{H}^1(\text{Cone}(\omega)) = E_{C,t} \oplus E_{C,t}^\perp$. Let $p_{C,t}, p_{C,t}^\perp$ denote the orthogonal projections from $\mathbf{H}^1(\text{Cone}(\omega))$ to $E_{C,t}$ and $E_{C,t}^\perp$, respectively, and decompose the operator $D_{C,t} = d_{C,t} + d_{C,t}^*$ via

$$D_{C,t,1} = p_{C,t} D_{C,t} p_{C,t}, \quad D_{C,t,2} = p_{C,t} D_{C,t} p_{C,t}^\perp, \quad D_{C,t,3} = p_{C,t}^\perp D_{C,t} p_{C,t}, \quad D_{C,t,4} = p_{C,t}^\perp D_{C,t} p_{C,t}^\perp.$$

We have the following results:

Theorem 2.8. *There exists a constant $t_0 > 0$ such that*

(i) *for any $t \geq t_0$, and $0 \leq u \leq 1$, the operator*

$$D_{C,t,u} = D_{C,t,1} + D_{C,t,4} + u(D_{C,t,2} + D_{C,t,3}) = D_{C,t} + (u - 1)(D_{C,t,2} + D_{C,t,3})$$

is Fredholm;

(ii) *the operator $D_{C,t,4} : E_{C,t}^\perp \cap \mathbf{H}^1(\text{Cone}(\omega)) \rightarrow E_{C,t}^\perp$ is invertible.*

To prove these, we need the following inequalities:

Lemma 2.9. *There exists a constant $t_1 > 0$ such that for $\sigma \in E_{C,t}^\perp \cap \mathbf{H}^1(\text{Cone}(\omega)), \sigma' \in E_{C,t}$ and $t \geq t_1$, we have*

$$\begin{aligned} \|D_{C,t,2}\sigma\|_0 &\leq \frac{C_1\|\sigma\|_0}{t}, \\ \|D_{C,t,3}\sigma'\|_0 &\leq \frac{C_1\|\sigma'\|_0}{t}, \end{aligned}$$

for some positive constant C_1 .

Proof. Note that $D_{C,t,3}$ is the adjoint of $D_{C,t,2}$, so if we prove the first bound, we obtain

the second. Since $\rho_{p,1}(t), \rho_{p,2}(t)$ are supported in V_p , we have

$$\begin{aligned}
D_{C,t,2} \sigma &= p_{C,t} D_{C,t} p_{C,t}^\perp \sigma = p_{C,t} D_{C,t} \sigma \\
&= \sum_{p \in \text{Crit}(f)} \rho_{p,i}(t) \int_{V_p} \langle \rho_{p,i}(t), D_{C,t} \sigma \rangle_x dVol \\
&= \sum_{p \in \text{Crit}(f), n_f(p)=k} \rho_{p,1}(t) \int_{V_p} \left\langle D_{C,t} \frac{\gamma_p}{\Phi_{p,1}(t)} \begin{pmatrix} \zeta_k \\ \iota_Z \zeta_k \end{pmatrix}, \sigma \right\rangle_x dVol \\
&\quad + \sum_{q \in \text{Crit}(f), n_f(q)=k-1} \rho_{q,2}(t) \int_{V_q} \left\langle D_{C,t} \frac{\gamma_q}{\Phi_{q,2}(t)} \begin{pmatrix} -\tau \wedge \zeta_{k-1} \\ \zeta_{k-1} \end{pmatrix}, \sigma \right\rangle_x dVol.
\end{aligned}$$

And note that γ_p is constant on $|\mathbf{x}| < a, |\mathbf{x}| > 2a$, so $D_{C,t} \rho_{p,1}(t) = 0 = D_{C,t} \rho_{p,2}(t)$ (as these are harmonic solutions multiplied by a constant) on $|\mathbf{x}| < a, |\mathbf{x}| > 2a$.

Now note that we chose

$$\Phi_{p,1}(t) = \sqrt{\int_{V_p} \gamma_p^2 \left\| \begin{pmatrix} \zeta_k \\ \iota_Z \zeta_k \end{pmatrix} \right\|_x^2 dVol}, \quad \Phi_{p,2}(t) = \sqrt{\int_{V_p} \gamma_p^2 \left\| \begin{pmatrix} -\tau \wedge \zeta_k \\ \zeta_k \end{pmatrix} \right\|_x^2 dVol},$$

so that $\rho_{p,1}(t) = \frac{\gamma_p}{\Phi_{p,1}(t)} \begin{pmatrix} \zeta_k \\ \iota_Z \zeta_k \end{pmatrix}$ and $\rho_{p,2}(t) = \frac{\gamma_p}{\Phi_{p,2}(t)} \begin{pmatrix} -\tau \wedge \zeta_k \\ \zeta_k \end{pmatrix}$ have unit norm. Thus note that, as $\zeta_k, -\tau \wedge \zeta_k, \iota_Z \zeta_k$ are either $e^{-t|x|^2/2} dx_I$ or $x_j e^{-t|x|^2/2} dx_J$ where $dx_J = \iota_{\partial_{n+j}} dx_I$ or $dx_{n+j} \wedge dx_I$. Thus, in the region $a < |\mathbf{x}| < 2a$, these are bounded above by $\int_{\mathbb{R}^n} (1 + 2a^n) e^{-t|x|^2/2} dVol = \max(1, (2a)^n) \left(\frac{2t}{\pi}\right)^{2n/2}$, and bounded below by

$$\int_{B_a} \max(1, 2a^n) e^{-t|x|^2/2} dVol = C_2 \left(\frac{2t}{\pi}\right)^{2n/2},$$

and thus $C_3 t^n \leq \frac{1}{\Phi_{p,i}(t)} \leq C_4 t^n$. We will now look at the integrals

$$\int_{V_p} \left\langle D_{C,t} \left(\frac{\gamma_p}{\Phi_{p,1}(t)} \begin{pmatrix} \zeta_k \\ \iota_Z \zeta_k \end{pmatrix} \right), \sigma \right\rangle dVol, \quad \int_{V_p} \left\langle D_{C,t} \left(\frac{\gamma_p}{\Phi_{p,2}(t)} \begin{pmatrix} -\tau \wedge \zeta_k \\ \zeta_k \end{pmatrix} \right), \sigma \right\rangle dVol.$$

Note that we can restrict these to $a < |x| < 2a$, where $\gamma_p \zeta_k, \Lambda \zeta_k, \zeta_{k-1}, -\omega \wedge \zeta_{k-1}$ are

bounded by $C_5 e^{-ta^2}$. Also note that

$$\begin{aligned}
[d_t + d_t^*](\gamma_p \psi) &= (d\gamma_p) \wedge \psi + \gamma_p d_t \psi + (d^* + t\nabla f)\gamma_p \psi \\
&= (d\gamma_p) \wedge \psi + \gamma_p d_t \psi + (\partial_i \gamma_p \psi_I) \iota_{\partial_i} dx_I + t\nabla f \gamma_p \psi_I dx_I \\
&= (d\gamma_p) \wedge \psi + \gamma_p d_t \psi + (\partial_i \gamma_p) \psi_I \iota_{\partial_i} dx_I + \gamma_p \partial_i \psi_I \iota_{\partial_i} dx_I + \gamma_p t \nabla f \psi_I dx_I \\
&= (d\gamma_p) \wedge \psi + \gamma_p d_t \psi + (\partial_i \gamma_p) \psi_I \iota_{\partial_i} dx_I + \gamma_p d^* \psi_I + t \gamma_p \iota_{\nabla f} \psi \\
&= (d\gamma_p) \wedge \psi + \gamma_p d_t \psi + (\partial_i \gamma_p) \psi_I \iota_{\partial_i} dx_I + \gamma_p d_t^* \psi.
\end{aligned}$$

Therefore, for $\psi = \zeta_k, \iota_Z \zeta_k, -\tau \wedge \zeta_k$, the terms we can get in $D_{C,t}$ are $\Lambda\psi, \omega \wedge \psi$, which are bounded by $\max(1, 2a)e^{-ta^2/2}$, and $(d\gamma_p) \wedge \psi, \gamma_p d_t \psi + \gamma_p d_t^* \psi, (\partial_i \gamma_p) \psi_I \iota_{\partial_i} dx_I$ which are bounded by $C_6 \max(1, 2a)e^{-ta^2/2}$ for some positive constant C_6 . Therefore, we have

$$\begin{aligned}
\int_{V_p} \left\langle D_{C,t} \left(\frac{\gamma_p}{\Phi_{p,1,t}} \begin{pmatrix} \zeta_k \\ \iota_Z \zeta_k \end{pmatrix} \right), \sigma \right\rangle_x dVol &\leq \|\sigma\|_0 \sqrt{\int_{a < |x| < 2a} \frac{C_7}{\Phi_{p,1,t}} e^{-ta^2/2} dVol} \\
&\leq C_8 t^n e^{-ta^2/2} \|\sigma\|_0 \leq \frac{C_9 \|\sigma\|_0}{t}.
\end{aligned}$$

and also

$$\begin{aligned}
\int_{V_p} \left\langle D_{C,t} \left(\frac{\gamma_p}{\Phi_{p,2,t}} \begin{pmatrix} -\tau \wedge \zeta_k \\ \zeta_k \end{pmatrix} \right), \sigma \right\rangle_x dVol &\leq \|\sigma\|_0 \sqrt{\int_{a < |x| < 2a} \frac{C_{10}}{\Phi_{p,2,t}} e^{-ta^2/2} dVol} \\
&\leq C_{11} t^n e^{-ta^2/2} \|\sigma\|_0 \leq \frac{C_{12} \|\sigma\|_0}{t}.
\end{aligned}$$

Therefore, for $C_1 = \max(C_{12}, C_8)$ the norm of $D_{C,t,2}\sigma$ satisfies

$$\|D_{C,t,2}\sigma\| \leq \left\| \rho_{p,i}(t) \frac{C_1 \|\sigma\|_0}{t} \right\| = \frac{C_1 \|\sigma\|_0}{t},$$

and we thus have our first inequality, and thus using properties of adjoints we have proved both inequalities. \square

Next, note that this implies that $D_{C,t,2}$ and $D_{C,t,3}$ are compact operators, and thus $D_{C,t,u} = D_{C,t} + (u-1)(D_{C,t,2} + D_{C,t,3})$ is a Fredholm operator plus a compact operator, hence also Fredholm. To show that the operator $D_{C,t,4} : E_{C,t}^\perp \cap \mathbf{H}^1(\text{Cone}(\omega)) \rightarrow E_{C,t}^\perp$ is invertible, we shall follow Bismut and Zhang [3] to show that there exists a con-

stant $t_2 > 0$ such that for $t \geq t_2$ and $\sigma \in E_{C,t}^\perp$,

$$\|D_{C,t,4}\sigma\|_0 \geq C_{13}\sqrt{t}\|\sigma\|_0. \quad (2.35)$$

To show this, we consider the inequality first in two special cases:

Case 1: $\text{supp}(\sigma) \subset V_p(4a)$ (the ball of radius $4a$).

Let E_p be a Euclidean space containing V_p , and $\mathbf{H}^0(E_p)$ be the Sobolev space corresponding to $\|\cdot\|_0$, the 0-th Sobolev norm on E_p . Define $\rho'_{p,1,t} = (\frac{t}{\pi})^{n/2} e^{-t|x|^2/2} \rho_{p,1,t}$ and $\rho'_{p,2,t} = (\frac{t}{\pi})^{n/2} e^{-t|x|^2} \rho_{p,2,t}$, and define $p'_{C,t}$ to be the projection onto the subspace of $\mathbf{H}^0(E_p)$ spanned by the $\rho'_{p,i,t}$. Since $\sigma \in E_{C,t}^\perp \cap \mathbf{H}^1(\text{Cone}(\omega))$, we have that $p_{C,t}$ projecting σ to $E_{C,t}$ is zero, i.e. $p_{C,t}\sigma = 0$. Accordingly, we have that

$$\begin{aligned} p'_{C,t}\sigma &= p'_{C,t}\sigma - p_{C,t}\sigma \\ &= \sum_{p \in \text{Crit}(f)} \rho'_{p,i,t} \int_{E_p} (1 - \gamma_p(|x|)) \left(\frac{t}{\pi}\right)^{n/2} e^{-t|x|^2/2} \langle \rho_{p,i,t}, \sigma \rangle_x dVol. \end{aligned}$$

As $\gamma_p = 1$ near p , and zero outside of the ball of radius $4a$, a similar calculation of the integral as was done for the proof of Lemma 2.9 shows

$$\|p'_{C,t}\sigma\|^2 \leq \frac{C_{14}}{\sqrt{t}} \|\sigma\|^2. \quad (2.36)$$

Next, note that $D_{C,t}\rho'_{p,i,t} = 0$, so $D_{C,t}p'_{C,t}\sigma = 0$, and as $\sigma - p'_{C,t}\sigma \in (E'_{C,t})^\perp$ we can apply inequalities 2.35 and 2.36 and get

$$\|D_{C,t}\sigma\|_0^2 = \|D_{C,t}(\sigma - p'_{C,t}\sigma)\|_0^2 \geq C_{13}t\|\sigma - p'_{C,t}\sigma\|_0^2 \geq C_{13}t\|\sigma\|_0^2 - C_{16}\sqrt{t}\|\sigma\|_0^2.$$

Thus, $\|D_{C,t}\sigma\|_0 \geq \frac{C_{17}\sqrt{t}}{2}\|\sigma\|_0$.

Case 2: $\text{supp}(\sigma) \subset M \setminus \bigcup_{p \in \text{Crit}(f)} V_p(2a)$ (and still $\sigma \in E_{C,t}^\perp \cap H^1(M)$).

To prove the estimate in this case, recall from (2.8) that

$$\begin{aligned}
D_{C,t}^2 &= \Delta_{C,t} = \begin{pmatrix} \Delta_t + \omega\Lambda & -d_t^{\Lambda^*} \\ -d_t^\Lambda & \Delta_t + \Lambda\omega \end{pmatrix} \\
&= \begin{pmatrix} \Delta + t(\mathcal{L}_{\nabla f} + \mathcal{L}_{\nabla f}^*) + t^2\|df\|_x^2 + \omega\Lambda & -d^* + t(\iota_{\nabla f}\omega - \omega\iota_{\nabla f}) \\ -d^\Lambda + t(\Lambda df - df\Lambda) & \Delta + t(\mathcal{L}_{\nabla f} + \mathcal{L}_{\nabla f}^*) + t^2\|df\|_x^2 + \Lambda\omega \end{pmatrix} \\
&= \begin{pmatrix} \Delta + \omega\Lambda & -d^{\Lambda^*} \\ -d^\Lambda & \Delta + \Lambda\omega \end{pmatrix} + t \begin{pmatrix} \mathcal{L}_{\nabla f} + \mathcal{L}_{\nabla f}^* & \iota_{\nabla f}\omega - \omega\iota_{\nabla f} \\ \Lambda df - df\Lambda & \mathcal{L}_{\nabla f} + \mathcal{L}_{\nabla f}^* \end{pmatrix} + t^2\|df\|_x^2 I. \\
&= \Delta_C + t \begin{pmatrix} \mathcal{L}_{\nabla f} + \mathcal{L}_{\nabla f}^* & \iota_{\nabla f}\omega - \omega\iota_{\nabla f} \\ \Lambda df - df\Lambda & \mathcal{L}_{\nabla f} + \mathcal{L}_{\nabla f}^* \end{pmatrix} + t^2\|df\|_x^2 I.
\end{aligned}$$

Since we are away from the zeroes of df , $\|df\|_x^2 \geq C_{17}$. Since $\text{supp}(\sigma)$ is away from the zeroes, we have that

$$\langle t^2\|df\|_x^2 \sigma, \sigma \rangle = \int_M t^2\|df\|_x^2 (\eta \wedge * \eta + \xi \wedge * \xi) \geq \int_M t^2 C_{17} (\eta \wedge * \eta + \xi \wedge * \xi) = t^2 C_{18} \|\sigma\|^2.$$

Also, note that $D_C^2 = d_C d_C^* + d_C^* d_C$ is a positive operator, hence $\langle D_C \sigma, \sigma \rangle \geq 0$. Note further that $\begin{pmatrix} \mathcal{L}_{\nabla f} + \mathcal{L}_{\nabla f}^* & \iota_{\nabla f}\omega - \omega\iota_{\nabla f} \\ \Lambda df - df\Lambda & \mathcal{L}_{\nabla f} + \mathcal{L}_{\nabla f}^* \end{pmatrix}$ is a zeroth order operator with an operator norm C_{19} , so

$$\begin{aligned}
\|D_{C,t}\sigma\|^2 &= \langle D_{C,t}\sigma, D_{C,t}\sigma \rangle = \langle D_{C,t}^2 \sigma, \sigma \rangle \\
&= \left\langle \left(D_C^2 + t \begin{pmatrix} \mathcal{L}_{\nabla f} + \mathcal{L}_{\nabla f}^* & \iota_{\nabla f}\omega - \omega\iota_{\nabla f} \\ \Lambda df - df\Lambda & \mathcal{L}_{\nabla f} + \mathcal{L}_{\nabla f}^* \end{pmatrix} + t^2\|df\|_x^2 I \right) \sigma, \sigma \right\rangle \\
&= \langle \Delta_C \sigma, \sigma \rangle + t \left\langle \begin{pmatrix} \mathcal{L}_{\nabla f} + \mathcal{L}_{\nabla f}^* & \iota_{\nabla f}\omega - \omega\iota_{\nabla f} \\ \Lambda df - df\Lambda & \mathcal{L}_{\nabla f} + \mathcal{L}_{\nabla f}^* \end{pmatrix} \sigma, \sigma \right\rangle + t^2 \langle \|df\|_x^2 \sigma, \sigma \rangle \\
&\geq (0 - C_{19}t + C_{18}t^2) \|\sigma\|^2,
\end{aligned}$$

from which we can conclude $\|D_{C,t}\sigma\| \geq C_{20}\sqrt{t}\|\sigma\|$.

With the two cases at hand, we now derive the inequality (2.35). First, we define the function $\tilde{\gamma}_p \in C^\infty(M)$ such that $\tilde{\gamma}_p(y) = \gamma_p(|y|/2)$ in V_p and $\tilde{\gamma}_p|_{M \setminus \cup V_p(4a)} = 0$. For $\sigma \in E_{C,t}^\perp \cap \mathbf{H}^1(M)$, one can see that $\tilde{\gamma}_p \sigma \in E_{C,t}^\perp \cap \mathbf{H}^1(M)$. Also, $\|D_{C,t}\sigma\| \geq \|D_{C,t}\sigma - \tilde{\gamma}_p D_{C,t}\sigma\| + \|\tilde{\gamma}_p D_{C,t}\sigma\|$.

Therefore, with Case 1 and 2, we deduce that there is a C_{21} such that for $t \geq t_1 + t_2$,

$$\begin{aligned}
\|D_{C,t}\sigma\|_0 &\geq \frac{1}{2}(\|(1 - \tilde{\gamma}_p)D_{C,t}\sigma\|_0 + \|\tilde{\gamma}_p D_{C,t}\sigma\|_0) \\
&\geq \frac{1}{2}(\|D_{C,t}(1 - \tilde{\gamma}_p)\sigma + [D_{C,t}, \gamma_p]\sigma\|_0 + \|D_{C,t}\tilde{\gamma}_p\sigma + [\gamma_p, D_{C,t}]\sigma\|_0) \\
&\geq \frac{\sqrt{t}}{2}(C_{20}\|(1 - \tilde{\gamma}_p)\sigma\|_0 + \sqrt{C_{17}}\|\tilde{\gamma}_p\sigma\|_0) - C_{21}\|\sigma\|_0 \\
&\geq \sqrt{t}C_{22}\|\sigma\|_0 - C_{21}\|\sigma\|_0
\end{aligned}$$

where $C_{22} = \min\{\sqrt{C_{17}}/2, C_{20}/2\}$. Thus, we have completed the proof of inequality (2.35) from which we conclude that the operator $D_{C,t,A} : E_{C,t}^\perp \cap \mathbf{H}^1(M) \rightarrow E_{C,t}^\perp$ is invertible when t is sufficiently large.

Lemma 2.10. *There exists a constant $C_{23} > 0$ such that for $\sigma \in \mathbf{H}^1(\text{Cone}(\omega))$ and any $t \geq t_3$, we have*

$$\|D_{C,t,1}\sigma\|_0 \leq \frac{C_{23}\|\sigma\|_0}{t}.$$

Proof. If we examine the integral, then using the $\rho_{p,i}$'s in (2.34) as our basis for $E_{C,t}$, we find

$$p_{C,t}\sigma = \sum_{p \in \text{Crit}(f)} \rho_{p,i}(t) \int_{V_p} \langle \rho_{p,i}(t), \sigma \rangle_x dVol.$$

And note that if we take $D_{C,t}\rho_{p,i}(t)$, then this is zero in the region $|x - p| < a$ and $|x - p| > 2a$, and from a similar argument to Lemma 2.9, $\|D_{C,t}\rho_{p,i}(t)\| \leq \frac{C_{12}}{t}$. Then using

that $\rho_{p,i}(t)$ has unit norm, we obtain

$$\begin{aligned}
\|p_{C,t}D_{C,t}p_{C,t}\sigma\|_0 &= \left\| \sum_{p \in \text{Crit}(f)} \rho_{p,i}(t) \int_{V_p} D_{C,t} \rho_{p,i}(t) \langle \rho_{p,i}(t), \sigma \rangle_x d\text{Vol} \right\|_0 \\
&\leq \sum_{p \in \text{Crit}(f)} \left\| \rho_{p,i}(t) \langle \rho_{p,i}(t), \sigma \rangle \int_{V_p} D_{C,t} \rho_{p,i}(t) d\text{Vol} \right\|_0 \\
&\leq \sum_{p \in \text{Crit}(f)} \left\| \langle \rho_{p,i}(t), \sigma \rangle \frac{C_{12}}{t} \right\|_0 \\
&\leq \sum_{p \in \text{Crit}(f)} \frac{C_{12}}{t} \|\rho_{p,i}(t)\|_0 \|\sigma\|_0 \\
&\leq \frac{C_{23}\|\sigma\|_0}{t}.
\end{aligned}$$

□

Definition 2.11. For any $b > 0$, let $E_{C,t}(b)$ denote the direct sum of eigenspaces of $D_{C,t}$ with eigenvalues in $[-b, b]$. Since $D_{C,t}$ is a self-adjoint linear operator, $E_{C,t}(b)$ is a finite-dimensional subspace of $\mathbf{H}^0(\text{Cone}(\omega))$.

Let $p_{C,t}(b)$ denote the projection operator from $\mathbf{H}^0(\text{Cone}(\omega))$ to $E_{C,t}(b)$.

Lemma 2.12. *There exists a $C_{24} > 0$ such that for $t \geq t_4$ and $\sigma \in E_{C,t}$,*

$$\|p_{C,t}(b)\sigma - \sigma\|_0 \leq \frac{C_{24}}{t} \|\sigma\|_0.$$

Proof. Let $\delta = \{\lambda \in \mathbb{C} : |\lambda| = C_\lambda\}$ be the counterclockwise oriented circle at radius C_λ . By Lemma 2.9 and Lemma 2.10, we have that for any $\lambda \in \delta$, $t \geq t_1 + t_2$, and $\sigma' \in \mathbf{H}^1(\text{Cone}(\omega))$,

$$\begin{aligned}
&\|(\lambda - D_{C,t})\sigma'\|_0 \\
&\geq \frac{1}{2} \left(\|\lambda p_{C,t}\sigma' - D_{t,1}p_{C,t}\sigma' - D_{t,2}p_{C,t}^\perp\sigma'\|_0 + \|\lambda p_T^\perp\sigma' - D_{t,3}p_{C,t}\sigma' - D_{t,4}p_{C,t}^\perp\sigma'\|_0 \right) \\
&\geq \frac{1}{2} \left(\left(C_\lambda - \frac{C_{12}}{t} - \frac{C_{23}}{t} \right) \|p_{C,t}\sigma'\|_0 + \left(C_{17}\sqrt{t} - C_\lambda - \frac{C_{23}}{t} \right) \|p_{C,t}^\perp\sigma'\|_0 \right).
\end{aligned}$$

By the above inequality, for $t_4 \geq t_1 + t_2$ and $C_{25} > 0$ such that for any $t > t_4$ and $\sigma' \in \mathbf{H}^1(\text{Cone}(\omega))$

$$\|(\lambda - D_{C,t})\sigma'\|_0 \geq C_{25}\|\sigma'\|_0.$$

Thus, for any $\lambda \in \delta$, $\lambda - D_{C,t} : \mathbf{H}^1(\text{Cone}(\omega)) \rightarrow \mathbf{H}^0(\text{Cone}(\omega))$ is invertible, so the resolvent $(\lambda - D_{C,t})^{-1}$ is well defined. By the basic spectral theory for operators, for $\sigma \in E_{C,t}$, one has

$$p_{C,t}(b)\sigma - \sigma = \frac{1}{2\pi\sqrt{-1}} \int_{\delta} ((\lambda - D_{C,t})^{-1} - \lambda^{-1})\sigma d\lambda.$$

Since with $p_{C,t}$ the projection to $E_{C,t}$, we have $p_{C,t}^\perp\sigma = 0$. Thus, using the inequality above, we have

$$\begin{aligned} ((\lambda - D_{C,t})^{-1} - \lambda^{-1})\sigma &= \lambda^{-1}(\lambda - D_{C,t})^{-1}D_{C,t}\sigma \\ &= \lambda^{-1}(\lambda - D_{C,t})^{-1}(D_{C,t,1}\sigma + D_{C,t,3}\sigma). \end{aligned}$$

One deduces by Lemma 2.9 and above we have

$$\begin{aligned} \|(\lambda - D_{C,t})^{-1}(D_{C,t,1}\sigma + D_{C,t,3}\sigma)\|_0 &\leq C_{13}^{-1}\|D_{t,1}\sigma + D_{t,3}\sigma\|_0 \\ &\leq C_{25}^{-1} \left(\frac{C_{12} + C_{23}}{t} \right) \|\sigma\|_0 \end{aligned}$$

and plugging this into the integral gives

$$\begin{aligned} \|p_{C,t}(b)\sigma - \sigma\|_0 &= \left\| \frac{1}{2\pi\sqrt{-1}} \int_{\delta} ((\lambda - D_{C,t})^{-1} - \lambda^{-1})\sigma d\lambda \right\|_0 \\ &\leq \frac{1}{2\pi} \int_{\delta} \|\lambda^{-1}(\lambda - D_{C,t})^{-1}(D_{C,t,1}\sigma + D_{C,t,3}\sigma)\|_0 d\lambda \\ &\leq \frac{C_{26}}{2\pi} \int_{\delta} \|C_{25}^{-1} \left(\frac{C_{12} + C_{23}}{t} \right) \|\sigma\|_0 d\lambda \\ &\leq \frac{C_{24}}{t} \|\sigma\|_0. \end{aligned}$$

□

Theorem 2.13. *Let $F_{C,t}^{[0,b]}$ be the space of all eigenforms of $\Delta_{C,t}$ with eigenvalues in $[0, b]$. Then for t large enough, $(F_{C,t}^{[0,b]}, d_{C,t})$ is a chain complex with $\dim (F_{C,t}^{[0,b]})_k = m_k + m_{k-1}$.*

Proof. By applying Lemma 2.12 to the $\rho_{p,i,t}$'s when t is large enough, $p_{C,t}(b)\rho_{p,i,t}$ will be

linearly independent. (If they are not linearly dependent, then $\sigma = \rho_{p,i,t}$ and $\sigma' = a\rho'_{p',i',t}$ would have $p_{C,t}(b)\sigma = p_{C,t}(b)\sigma'$, but then by Lemma 2.12, we would have

$$\frac{C_1}{t}\|\sigma - \sigma'\|_0 \geq \|p_{C,t}(b)\sigma - p_{C,t}(b)\sigma' - (\sigma - \sigma')\|_0 = \|\sigma - \sigma'\|_0$$

which is a contradiction.) Thus for $t \geq t_5$, we have $\dim(E_{C,t}(b)) \geq \dim(E_{C,t})$. Now assume for the purposes of contradiction that $\dim(E_{C,t}(b)) \geq \dim(E_{C,t})$. Then there is a nonzero $\sigma \in E_{C,t}(b)$ that is orthogonal to $p_{C,t}(b)E_{C,t}$, or $\langle \sigma, p_{C,t}\rho_{p,i}(t) \rangle_{\mathbf{H}^0(\text{Cone}(\omega))} = 0$ for any $\rho_{p,i}$. Then from Lemma 2.9 and Case 1, we have that

$$\begin{aligned} p_{C,t}\sigma &= \sum_{p \in \text{Crit}(f)} \langle \sigma, \rho_{p,i}(t) \rangle \rho_{p,i}(t) \\ &= \sum_{p \in \text{Crit}(f)} \langle \sigma, \rho_{p,i}(t) \rangle \rho_{p,i}(t) - \sum_{p \in \text{Crit}(f)} \langle \sigma, p_{C,t}\rho_{p,i}(t) \rangle p_{C,t}(b)\rho_{p,i}(t) \\ &= \sum_{p \in \text{Crit}(f)} \langle \sigma, \rho_{p,i}(t) \rangle (\rho_{p,i}(t) - p_{C,t}(b)\rho_{p,i}(t)) + \sum_{p \in \text{Crit}(f)} \langle \sigma, \rho_{p,i}(t) - p_{C,t}\rho_{p,i}(t) \rangle p_{C,t}(b)\rho_{p,i}(t). \end{aligned}$$

By Lemma 2.9, there exists a $C_{12} > 0$ so when $t \geq t_5$ $\|p_{C,t}\sigma\|_0 \leq \frac{C_{12}}{t}\|\sigma\|_0$, and thus

$$\|p_{C,t}^\perp\sigma\|_0 = (\|\sigma\|_0 - \|p_{C,t}\sigma\|_0) \geq \|\sigma\|_0 \geq C_{25}\|\sigma\|_0.$$

Using this, (2.35), Lemma 2.12, and when $t > 0$ is large enough, we have

$$\begin{aligned} C_{25}C_{20}\sqrt{t}\|\sigma\|_0 &\leq C_{20}\sqrt{t}\|p_{C,t}^\perp\sigma\|_0 \\ &\leq \|D_{C,t}p_{C,t}^\perp\sigma\|_0 \\ &= \|D_{C,t}\sigma - D_{C,t}p_{C,t}\sigma\|_0 \\ &= \|D_{C,t}\sigma - D_{C,t,1}\sigma - D_{C,t,3}\sigma\|_0 \\ &\leq \|D_{C,t}\sigma\|_0 + \|D_{C,t,1}\sigma\|_0 + \|D_{C,t,3}\sigma\|_0 \\ &\leq \|D_{C,t}\sigma\|_0 + \frac{C_{12} + C_{23}}{t}\|\sigma\|_0, \end{aligned}$$

from which one gets $\|D_{C,t}\sigma\|_0 \geq C_{25}C_{20}\sqrt{t}\|\sigma\|_0 - \frac{C_{12} + C_{23}}{t}\|\sigma\|_0$ which contradicts that $\sigma \in E_{C,t}(b)$ is an eigenspace of $D_{C,t}$ for t large enough. Thus, one has

$$\dim(E_{C,t}(b)) = \dim E_{C,t} = \sum_k m_k + m_{k-1} = 2 \sum_k m_k.$$

Moreover, $E_{C,t}$ is generated by $p_{C,t}(b) \rho_{p,i}(t)$. □

Now to prove Theorem 2.13, for any integer k , such that $0 \leq k \leq 2n+1$, let Q_i denote the projection from $\mathbf{H}^0(\text{Cone}(\omega))$ onto the L^2 completion of $\text{Cone}^k(\omega)$. Since $\Delta_{C,t}$ preserves the \mathbb{Z} grading of $\Omega^*(M)$, for any eigenvector σ of $D_{C,t}$ associated with an eigenvalue $\mu \in [-b, b]$

$$\Delta_{C,t} Q_k \sigma = Q_k \Delta_{C,t} \sigma = Q_k \mu^2 \sigma = \mu^2 Q_k \sigma.$$

That is, $Q_k \sigma$ is an eigenform of $\Delta_{C,t}$ with eigenvalue μ^2 . We thus need to show that $\dim Q_k E_{C,t}(b) = m_k + m_{k-1}$. To prove this, note that by Lemma 2.12,

$$\|Q_{n_f(p)} p_{C,t}(b) \rho_{p,i}(t) - \rho_{p,i}(t)\|_0 \leq \frac{C_{24}}{t}.$$

Thus, for t sufficiently large, the cone forms $Q_{n_f(p)} p_{C,t}(b) \rho_{p,i}(t)$ are linearly independent. Therefore, for each k , we have

$$\dim Q_k E_{C,t}(b) \geq m_k + m_{k-1}.$$

However, we also have (as every element in $\mathbf{H}^0(\text{Cone}(\omega))$ is a linear combination of $2n+1$ form)

$$\sum_{k=0}^{2n+1} \dim Q_k E_{C,t}(b) = \sum_{k=0}^{2n+1} \dim E_{C,t}(b) = \sum_k m_k + m_{k-1} = 2 \sum_k m_k.$$

From this and $\dim Q_k E_{C,t}(b) \geq m_k + m_{k-1}$, we obtain

$$\dim Q_k E_{C,t}(b) = m_k + m_{k-1}.$$

3 Relation between the cone complex and the cone Morse complex

For $(F_{C,t}^{[0,1]})_k$, the space of all eigenforms of $\Delta_{C,t}$ in $\text{Cone}^k(\omega)$ with eigenvalues in $[0, 1]$, we point out that $d_{C,t} : (F_{C,t}^{[0,1]})_k \rightarrow (F_{C,t}^{[0,1]})_{k+1}$, since $[\Delta_{C,t}, d_{C,t}] = 0$. Hence, $((F_{C,t}^{[0,1]})_\bullet, d_{C,t})$ is a cochain complex with cohomology, $H^k(F_{C,t}^{[0,1]}) \cong H^k(\text{Cone}(\omega))$. By Theorem 2.13, when t is sufficiently large, we have

$$\dim (F_{C,t}^{[0,1]})_k = m_k + m_{k-1}, \quad k = 0, 1, \dots, 2n+1. \quad (3.1)$$

which is precisely equal to the dimension of the cone Morse complex $\text{Cone}^k(c(\omega)) = C^k(M, f) \oplus C^{k-1}(M, f)$ defined in Definition 1.6. As an immediate corollary, we have the bound

$$b_k^\omega = \dim H^k(\text{Cone}(\omega)) \leq m_k + m_{k-1}, \quad k = 0, 1, \dots, 2n + 1.$$

However, this inequality does not give a sharp bound. Below, we shall proceed to prove the isomorphism of $H^k(\text{Cone}(\omega)) \cong H^k(\text{Cone}(c(\omega)))$ Theorem 1.3, which will allow us to derive the sharp Morse inequality bounds described in Theorem 1.4.

3.1 Quasi-isomorphism between $\text{Cone}(\omega)$ and $\text{Cone}(c(\omega))$

In this subsection, we will prove Theorem 1.3 by showing $H^k(\text{Cone}(\omega)) \cong H^k(\text{Cone}(c(\omega)))$. To do so, let us first briefly review the relationship between de Rham cohomology, $H_{dR}^k(M)$, and Morse cohomology, $H_{C(f)}^k(M)$. Recall that there is a map \mathcal{P} that links the de Rham complex with the Morse complex [3].

Definition 3.1. Define the map $\mathcal{P} : \Omega^k(M) \rightarrow C^k(M, f)$ by

$$\mathcal{P}\phi = \sum_{p_k \in \text{Crit}(f)} \left(\int_{U_{p_k}} \phi \right) p_k$$

where $\phi \in \Omega^k(M)$ and U_p is the unstable submanifold consisting of gradient flow lines moving away from p .

Importantly, \mathcal{P} is a chain map and induces an isomorphism on cohomology.

Theorem 3.2. ([3, Theorem 2.9]) *The map $\mathcal{P} : \Omega^k(M) \rightarrow C^k(M, f)$ is a chain map, i.e.*

$$\partial \mathcal{P} = \mathcal{P} d, \tag{3.2}$$

and moreover,

$$[\mathcal{P}] : H_{dR}^k(M) \rightarrow H_{C(f)}^k(M) \text{ is an isomorphism.} \tag{3.3}$$

The analytical Witten deformation proof of this theorem can be described by the fol-

lowing diagram:

$$\begin{array}{ccccc}
(\Omega^\bullet(M), d) & \xleftarrow{e^{tf}} & (\Omega^\bullet(M), d_t) & \xleftarrow{\iota_t^{[0,1]}} & ((F_t^{[0,1]})_\bullet, d_t) \\
\mathcal{P} \downarrow & & \swarrow \mathcal{P}_t & & \swarrow \mathcal{P}_t|_{F_t^{[0,1]}} \\
(C^\bullet(M, f), \partial) & & & &
\end{array} \tag{3.4}$$

where $(\Omega^\bullet(M), d_t)$ is the Witten-deformed de Rham complex, and similar to $((F_{C,t}^{[0,1]})_\bullet, d_{C,t})$, $((F_t^{[0,1]})_\bullet, d_t)$ is the cochain complex consisting of eigenforms of Δ_t in $\Omega(M)$ with eigenvalues in $[0, 1]$. In the diagram, $\iota_t^{[0,1]}$ is the inclusion map, and also, the \mathcal{P} map induces the map

$$\mathcal{P}_t = \mathcal{P} e^{tf} : (\Omega^k(M), d_t) \rightarrow (C^k(M, f), \partial). \tag{3.5}$$

The proof of Theorem 3.2 involves showing that for t large enough, $\mathcal{P}_t|_{F_t^{[0,1]}} : ((F_t^{[0,1]})_\bullet, d_t) \rightarrow (C^\bullet(M, f), \partial)$ is a cochain isomorphism [20]. This implies that the vertical map \mathcal{P} gives an isomorphism on cohomology since the cohomologies of $\Omega^\bullet(M)$ and $F_t^{[0,1]}$ are always identical, regardless of the value of t .

In considering $\text{Cone}(c(\omega))$, let us first recall the definition of the map $c(\omega) : C^k(M, f) \rightarrow C^{k+2}(M, f)$ when acting on a critical point $p \in \text{Crit}(f)$ with index $n_f(p) = k$:

$$c(\omega) p = \sum_{q \in \text{Crit}(f)} \left(\int_{\mathcal{M}(q,p)} \omega \right) q, \tag{3.6}$$

where the sum is over critical points q with index $n_f(q) = n_f(p) + 2$ and $\mathcal{M}(q, p)$ is the submanifold of flow lines from q to p . It was shown in [1, Section 3.5] and [18, Lemma 4] that

$$[\mathcal{P}][\omega] = [c(\omega)][\mathcal{P}], \tag{3.7}$$

that is, they are cohomologous as maps from $H_{dR}^k(M)$ to $H_{C(f)}^{k+2}(M)$.

In the following, we assume that t is sufficiently large such that $\mathcal{P}_t = \mathcal{P} e^{tf}$ is an isomorphism between $(F_t^{[0,1]})_k$ and $C^k(M, f)$. By (3.7), we have that

$$[\mathcal{P}_t][\omega] = [c(\omega)][\mathcal{P}_t]$$

on cohomology. This motivates us to introduce an $\omega \wedge$ type map on $F_t^{[0,1]}$.

Definition 3.3. For t sufficiently large, we define the map $\tilde{\omega}_t : (F_t^{[0,1]})_k \rightarrow (F_t^{[0,1]})_{k+2}$ by

$$\tilde{\omega}_t = \mathcal{P}_t^{-1} c(\omega) \mathcal{P}_t.$$

Since $d_t = e^{-tf} d e^{tf}$, we can extend $\mathcal{P} d = \partial \mathcal{P}$ of (3.2) to

$$\mathcal{P}_t d_t = \partial \mathcal{P}_t. \quad (3.8)$$

Applying \mathcal{P}_t^{-1} on both the left and the right of (3.8) gives $d_t \mathcal{P}_t^{-1} = \mathcal{P}_t^{-1} \partial$ acting on $C(M, f)$. Using the commutativity of $c(\omega)$ with ∂ , i.e. $\partial c(\omega) = c(\omega) \partial$ in (1.8), it is straightforward to check that $\tilde{\omega}$ is a chain map:

$$\begin{aligned} \tilde{\omega}_t d_t &= \mathcal{P}_t^{-1} c(\omega) \mathcal{P}_t d_t \\ &= \mathcal{P}_t^{-1} c(\omega) \partial \mathcal{P}_t \\ &= \mathcal{P}_t^{-1} \partial c(\omega) \mathcal{P}_t \\ &= d_t \mathcal{P}_t^{-1} c(\omega) \mathcal{P}_t = d_t \tilde{\omega}_t. \end{aligned}$$

The induced map on the cohomology

$$[\tilde{\omega}_t] = [\mathcal{P}_t^{-1}] [c(\omega)] [\mathcal{P}_t] = [\mathcal{P}_t^{-1}] [\mathcal{P}] [\omega \wedge] [\mathcal{P}]^{-1} [\mathcal{P}_t], \quad (3.9)$$

is thus conjugate to the wedge product map $\omega \wedge$. With $\tilde{\omega}_t$, we can use it to define the following cone complex:

$$\text{Cone}^k(\tilde{\omega}_t) = (F_t^{[0,1]})_k \oplus (F_t^{[0,1]})_{k-1}, \quad \tilde{d}_{C,t} = \begin{pmatrix} d_t & \tilde{\omega}_t \\ 0 & -d_t \end{pmatrix}. \quad (3.10)$$

The cohomology of this cone Morse complex $(\text{Cone}^\bullet(\tilde{\omega}_t), \tilde{d}_t)$ can be expressed in terms of the cokernels and kernels of the $\tilde{\omega}_t$ map in the following way:

$$\begin{aligned} H^k(\text{Cone}(\tilde{\omega}_t))(M) &\cong \text{coker} \left[[\tilde{\omega}_t] : H^{k-2}(F_t^{[0,1]}) \rightarrow H^k(F_t^{[0,1]}) \right] \\ &\quad \oplus \text{ker} \left[[\tilde{\omega}_t] : H^{k-1}(F_t^{[0,1]}) \rightarrow H^{k+1}(F_t^{[0,1]}) \right]. \end{aligned} \quad (3.11)$$

This relation follows from the fact that the cone complex with elements $\text{Cone}^k(\tilde{\omega}_t)$ sits in

a short exact sequence of chain complexes

$$0 \longrightarrow ((F_t^{[0,1]})_k, \tilde{d}_t) \xrightarrow{\iota} (\text{Cone}^k(\tilde{\omega}_t), \tilde{d}_{C,t}) \xrightarrow{\pi} ((F_t^{[0,1]})_{k-1}, -\tilde{d}_t) \longrightarrow 0 \quad (3.12)$$

where ι is the inclusion map and π is the projection onto the second component. The short exact sequence implies the following long exact sequence of cohomologies

$$\begin{array}{c} \dots \longrightarrow H^{k-2}(F_t^{[0,1]}) \xrightarrow{[\tilde{\omega}_t]} H^k(F_t^{[0,1]}) \xrightarrow{[\iota]} H^k(\text{Cone}(\tilde{\omega}_t)) \\ \left. \begin{array}{l} \xrightarrow{[\pi]} \\ \xrightarrow{[\tilde{\omega}_t]} \end{array} \right\} \hspace{10em} (3.13) \\ \xrightarrow{[\pi]} H^{k-1}(F_t^{[0,1]}) \xrightarrow{[\tilde{\omega}_t]} H^{k+1}(F_t^{[0,1]}) \longrightarrow \dots \end{array}$$

which implies (3.11).

We now point out two important properties of $H^k(\text{Cone}(\tilde{\omega}_t))$ when t is sufficiently large such that \mathcal{P}_t is an isomorphism.

- (i) $H^k(\text{Cone}(\tilde{\omega}_t)) \cong H^k(\text{Cone}(c(\omega)))$. By construction, for t sufficiently large, the complex $(\text{Cone}^\bullet(\tilde{\omega}_t), \tilde{d}_{C,t})$ is isomorphic to the cone Morse complex $(\text{Cone}^\bullet(c(\omega)), \partial_C)$. Hence, their cohomologies must be isomorphic.
- (ii) $H^k(\text{Cone}(\tilde{\omega}_t)) \cong H^k(\text{Cone}(\omega))$. Being both cone cohomologies, both $H^k(\text{Cone}(\tilde{\omega}_t))$ and $H^k(\text{Cone}(\omega))$ can be expressed in terms of the cokernels and kernels, of the $[\tilde{\omega}_t]$ map (1.3) and the $[\omega]$ map (3.11), respectively. Moreover, since $H^k(F_t^{[0,1]}) \cong H_{dR}^k(M)$ and also $[\tilde{\omega}_t]$ and $[\omega]$ have equivalent action on the cohomology level by (3.9), the two cohomologies are isomorphic.

Together, they imply the desired isomorphism that $H^k(\text{Cone}(c(\omega))) \cong H^k(\text{Cone}(\omega)) \cong PH^k(M, \omega)$ for $k = 0, 1, \dots, 2n + 1$. And this proves Theorem 1.3.

3.2 Cone Morse inequalities

Having proved $H(\text{Cone}(\omega)) \cong H(\text{Cone}(c(\omega)))$, we now proceed to derive Morse-type bounds for $b_k^\omega = \dim PH^k(M, \omega) = \dim H^k(\text{Cone}(\omega))$.

To do so, we note that the cohomology of the cone Morse complex $(\text{Cone}^\bullet(c(\omega)), \partial_C)$, like any cone cohomology, can be expressed in terms of cokernels and kernels of the $c(\omega)$

map:

$$\begin{aligned} H^k(\text{Cone}(c(\omega)))(M) &\cong \text{coker} \left[[c(\omega)] : H_{C(f)}^{k-2}(M) \rightarrow H_{C(f)}^k(M) \right] \\ &\oplus \ker \left[[c(\omega)] : H_{C(f)}^{k-1}(M) \rightarrow H_{C(f)}^{k+1}(M) \right] \end{aligned} \quad (3.14)$$

where $H_{C(f)}(M)$ is the cohomology of the standard Morse cochain complex $(C^\bullet(M, f), \partial)$. This relation can be derived similarly as that for $H^k(\text{Cone}(\tilde{\omega}_t))$ in (3.11) by means of short exact sequence of chain complexes (3.12) resulting in a long exact sequence of cohomologies (3.13).

Since the dimensions of the cohomology of the Morse complex are given by the Betti numbers, i.e. $b_k = \dim H_{C(f)}^k(M)$, it follows from the isomorphism $H^k(\text{Cone}(\omega)) \cong H^k(\text{Cone}(c(\omega)))$ and (3.14) that

$$b_k^\omega = \dim H^k(\text{Cone}(\omega)) = b_k - r_{k-2} + b_{k-1} - r_{k-1} \quad (3.15)$$

where

$$\begin{aligned} r_k &= \text{rank} \left([c(\omega)] : H_{C(f)}^k(M) \rightarrow H_{C(f)}^{k+2}(M) \right) \\ &= \text{rank} \left([\omega] : H_{dR}^k(M) \rightarrow H_{dR}^{k+2}(M) \right) \end{aligned} \quad (3.16)$$

is the rank of the $c(\omega)$ map on $H_{C(f)}^k(M)$, or equivalently, by (3.3) and (3.7), the rank of the ω map on $H_{dR}^k(M)$ as expressed in the second line of (3.16).

We can now bound the b_k^ω 's in terms of the critical points of the Morse function f using the following two properties. First, from the standard Morse theory, we have that the Betti numbers are bounded by

$$b_k \leq m_k, \quad k = 0, 1, \dots, 2n, \quad (3.17)$$

where m_k is the number index k critical points of f . Second, we introduce

$$v_k = \text{rank} \left(c(\omega) : C^k(M, f) \rightarrow C^{k+2}(M, f) \right) \quad (3.18)$$

which is the map of $c(\omega)$ on $C^k(M, f)$. Since r_k is the rank of the map on cohomology while v_k is the rank of the map on cochain, with both the cohomology $H_{C(f)}^k(M)$ and cochain

$C^k(M, f)$ generated by critical points, it is evident that

$$r_k \leq v_k, \quad k = 0, 1, \dots, 2n. \quad (3.19)$$

Now applying (3.17) and (3.19) to the expression for b_k^ω in (3.15), we immediately obtain the weak cone Morse inequalities

$$b_k^\omega \leq m_k - v_{k-2} + m_{k-1} - v_{k-1}. \quad (3.20)$$

Furthermore, we note that the alternating sum of b_k^ω 's results in the following

$$\begin{aligned} \sum_{i=0}^k (-1)^{k-i} b_i^\omega &= \sum_{i=0}^k (-1)^{k-i} (b_i - r_{i-2} + b_{i-1} - r_{i-1}) \\ &= b_k - r_{k-1} \leq m_k - v_{k-1}, \end{aligned} \quad (3.21)$$

which gives the strong cone Morse inequalities.

Finally, let us consider the case when f is a perfect Morse function. By definition, perfect Morse implies $b_k = m_k$ for all values of k . This means that $\dim H_{C(f)}^k = \dim C^k(M, f)$, and in particular, the Morse differential ∂ acts by zero on all generators of $C^*(M, f)$. Clearly then, when f is perfect, we have both $b_k = m_k$ and also $r_k = v_k$. Hence, both the weak cone Morse inequalities of (3.20) and the strong Morse inequalities of (3.21) would become equalities when f is a perfect Morse function.

Altogether, the weak and strong cone Morse inequalities and that they become equalities when f is perfect are the statements of Theorem 1.4. We have thus completed the proof of Theorem 1.4.

4 Examples

In this section, we will consider on certain symplectic manifolds the cone Morse complex and check the cone Morse inequalities derived in the previous section:

$$(weak) \quad b_k^\omega \leq m_k - v_{k-2} + m_{k-1} - v_{k-1}, \quad (4.1)$$

$$(strong) \quad \sum_{i=0}^k (-1)^{k-i} b_i^\omega \leq m_k - v_{k-1}, \quad (4.2)$$

for $k = 0, 1, \dots, 2n + 1$. In the first examples that we shall consider, the symplectic manifolds are Kähler. Due to the hard Lefschetz property, the wedge product map $[\omega^j] : H_{dR}^{n-j}(M) \rightarrow H_{dR}^{n+j}(M)$ for $j = 1, 2, \dots, n$, is an isomorphism. This implies in particular for $r_k = \text{rank } [\omega]_{H_{dR}^k(M)}$, that $r_k = \min(b_k, b_{k+2})$. It thus follows from the relation [12, 14]

$$H^k(\text{Cone}(\omega)) \cong \text{coker} \left[\omega : H_{dR}^{k-2}(M) \rightarrow H_{dR}^k(M) \right] \oplus \ker \left[\omega : H_{dR}^{k-1}(M) \rightarrow H_{dR}^{k+1}(M) \right] \quad (4.3)$$

that

$$b_k^\omega = \dim H^k(\text{Cone}(\omega)) = \begin{cases} b_k - b_{k-2} & 0 \leq k \leq n, \\ b_{k-1} - b_{k+1} & n + 1 \leq k \leq 2n + 1, \end{cases} \quad (4.4)$$

which are determined solely by the Betti numbers and do not vary with the symplectic structure. This is special to Kähler symplectic manifolds, as generally, the b_k^ω 's can vary with the class $[\omega] \in H_{dR}^2(M)$ (for explicit examples, see [15, 17]).

Remark 4.1. In the special case where the Morse function f and Riemannian metric g are chosen such that $c(\omega)^k : C^{n-k}(M, f) \rightarrow C^{n+k}(M, f)$ is bijective, mirroring the Lefschetz property but on the level of the Morse cochains, then $v_k = \text{rank } c(\omega) = \min(m_k, m_{k+2})$. The weak cone Morse inequalities would then also be analogous to (4.4)

$$b_k^\omega \leq \begin{cases} m_k - m_{k-2} & 0 \leq k \leq n, \\ m_{k-1} - m_{k+1} & n + 1 \leq k \leq 2n + 1. \end{cases} \quad (4.5)$$

While this does not hold generally, it always occurs when using a perfect Morse function manifolds on Kähler manifolds, which is the setting of our two Kähler examples below.

Example 4.2. Consider $(\mathbb{C}\mathbb{P}^n, \omega_{FS} = \frac{i}{2} \partial \bar{\partial} \log |z_i|^2)$, the complex n -dimensional projective space equipped with the Fubini-Study metric as the Kähler structure. It is well-known to have a perfect Morse function that can be expressed as

$$f([z_0, \dots, z_n]) = \frac{\sum \lambda_i |z_i|^2}{\sum |z_i|^2}$$

such that $\lambda_i \neq \lambda_j$ for $i \neq j$. If we consider $\lambda_0 < \lambda_1 < \dots < \lambda_n$, then the critical points are $p_{2j} = [0 \dots : 1 : \dots : 0]$ with 1 only in the j -th position and index $n_f(p_{2j}) = 2j$. Because the

index of all the critical points are even, the Morse differential $\partial : C^k(M, f) \rightarrow C^{k+1}(M, f)$ necessarily vanishes for all $k = 0, \dots, 2n$. Thus, f is no doubt a perfect Morse function and

$$m_k = b_k(\mathbb{C}\mathbb{P}^n) = \begin{cases} 1 & 0 \leq k \leq 2n, k \text{ even}, \\ 0 & \text{otherwise}. \end{cases} \quad (4.6)$$

Regarding the cone Morse differential $\partial_C = \begin{pmatrix} \partial & c(\omega_{FS}) \\ 0 & -\partial \end{pmatrix}$, it has a non-zero component coming from the $c(\omega_{FS})$ map. We note that

$$\mathcal{M}(p_{2j}, p_{2j+2}) = \{[0 : \dots : z_j : z_{j+1} : \dots : 0] : (z_j, z_{j+1}) \in \mathbb{C}^2 \setminus \{0\}\}$$

is isomorphic to $\mathbb{C}\mathbb{P}^1$. Therefore,

$$c(\omega_{FS})p_{2j} = \left(\int_{\mathbb{C}\mathbb{P}^1} \omega_{FS} \right) p_{2j+2} = \pi p_{2j+2},$$

and thus,

$$v_{2j} = \text{rank } c(\omega_{FS})|_{C^{2j}(f)} = 1, \quad j = 0, \dots, n-1, \quad (4.7)$$

which is the same as $r_{2j} = \text{rank } \omega|_{H^{2j}(\mathbb{C}\mathbb{P}^n)}$. The cone Morse complex and cohomology can then be easily computed and we find

$$b_k^\omega = \begin{cases} 1 & k = 0, 2n+1, \\ 0 & \text{otherwise}, \end{cases} \quad (4.8)$$

which agrees exactly with the expectation from (4.4).

The cone Morse inequalities (4.1)-(4.2) can similarly be straightforwardly checked using (4.6)-(4.7), and they are, in fact, equalities, as would be expected for a perfect Morse function.

Example 4.3. Consider $(T^4 = \mathbb{R}^4/\mathbb{Z}^4, \omega = dx_1 \wedge dx_2 + dx_3 \wedge dx_4)$, the four-torus described using Euclidean coordinates, x_i with identification $x_i \sim x_i + 1$ for $i = 1, 2, 3, 4$. For this example, we will compute the $\text{Cone}(c(\omega))$ complex with respect to the flat metric,

$g = \sum dx_i^2$, and the Morse function is taken to be

$$f = 2 - \frac{1}{2} \sum_{i=1}^4 \cos(2\pi x_i). \quad (4.9)$$

This Morse function has several desirable properties that are straightforward to prove:

- (i) the non-degenerate critical points are located at $x_i = [0]$ or $x_i = [\frac{1}{2}]$ and have Morse index equal to the number of coordinates which are equal to $[\frac{1}{2}]$;
- (ii) the number of critical points of index k , $m_k = b_k(T^4)$ for all k . Hence, f is perfect and the Morse differential ∂ acts by zero;
- (iii) the pair (f, g) satisfies Smale transversality.

Because of (ii), the ∂_C map reduces to the $c(\omega)$ map. Hence, we are interested in pairs of critical points whose indices differ by two, e.g. q_{k+1} has two more $[\frac{1}{2}]$ coordinates than q_{k-1} . Also, note that $\mathcal{M}(q_{k+1}, q_{k-1})$ will be a two-dimensional face with two of the coordinates fixed and two coordinates spanning the entire coordinate interval $[0, 1]$ when we take the closure. In Table 1, we give the cohomologies of $H(\text{Cone}(c(\omega)))$ and $H(\text{Cone}(\omega))$. We use a multi-index notation of $I = \{i_1 \dots i_j\}$ in increasing order such that $dx_I = dx_{i_1} \wedge \dots \wedge dx_{i_j}$, q_0 denotes the index 0 point, and q_I denotes the point with $\frac{1}{2}$ in entry i_1, \dots, i_j , i.e. $q_{13} = q_{[\frac{1}{2}, 0, \frac{1}{2}, 0]}$. The orientation of the submanifolds are chosen such that $\mathcal{P}dx_I = q_I$. (c.f. Definition 3.1.)

Notice that $c(\omega)q_I$ only picks out critical points that have two coordinates of q_I changed from $[0]$ to $[\frac{1}{2}]$ in either the 1-2 or 3-4 directions. Thus, we find that

$$\begin{aligned} c(\omega)q_0 &= q_{12} + q_{34}, & c(\omega)q_{12} &= q_{1234}, & c(\omega)q_{34} &= q_{1234}, \\ c(\omega)q_1 &= q_{134}, & c(\omega)q_2 &= q_{234}, & c(\omega)q_3 &= q_{123}, & c(\omega)q_4 &= q_{124}, \end{aligned}$$

with all other critical points mapped to zero when acted upon by $c(\omega)$.

It is straightforward to see from above that $v_k = r_k$ and that cone Morse inequalities give the equalities $b_k^\omega = m_k - v_{k-2} - m_{k-1} - v_{k-1}$ for $0 \leq k \leq 5$. This is as expected with f in (4.9) being a perfect Morse function.

Next, we consider a non-Kähler symplectic manifold where the hard Lefschetz property does not hold.

k	0	1	2
$H^k(\text{Cone}(\omega))$	$\begin{pmatrix} 1 \\ 0 \end{pmatrix}$	$\begin{pmatrix} dx_1 \\ 0 \end{pmatrix}, \begin{pmatrix} dx_2 \\ 0 \end{pmatrix}$ $\begin{pmatrix} dx_3 \\ 0 \end{pmatrix}, \begin{pmatrix} dx_4 \\ 0 \end{pmatrix}$	$\begin{pmatrix} dx_{13} \\ 0 \end{pmatrix}, \begin{pmatrix} dx_{14} \\ 0 \end{pmatrix}, \begin{pmatrix} dx_{23} \\ 0 \end{pmatrix},$ $\begin{pmatrix} dx_{24} \\ 0 \end{pmatrix}, \begin{pmatrix} dx_{12} - dx_{34} \\ 0 \end{pmatrix}$
$H^k(\text{Cone}(c(\omega)))$	$\begin{pmatrix} q_0 \\ 0 \end{pmatrix}$	$\begin{pmatrix} q_1 \\ 0 \end{pmatrix}, \begin{pmatrix} q_2 \\ 0 \end{pmatrix}$ $\begin{pmatrix} q_3 \\ 0 \end{pmatrix}, \begin{pmatrix} q_4 \\ 0 \end{pmatrix}$	$\begin{pmatrix} q_{13} \\ 0 \end{pmatrix}, \begin{pmatrix} q_{14} \\ 0 \end{pmatrix}, \begin{pmatrix} q_{23} \\ 0 \end{pmatrix},$ $\begin{pmatrix} q_{24} \\ 0 \end{pmatrix}, \begin{pmatrix} q_{12} - q_{34} \\ 0 \end{pmatrix}$
k	3	4	5
$H^k(\text{Cone}(\omega))$	$\begin{pmatrix} dx_{123} \\ 0 \end{pmatrix}, \begin{pmatrix} dx_{124} \\ 0 \end{pmatrix}, \begin{pmatrix} dx_{234} \\ 0 \end{pmatrix},$ $\begin{pmatrix} dx_{234} \\ 0 \end{pmatrix}, \begin{pmatrix} 0 \\ dx_{12} - dx_{34} \end{pmatrix}$	$\begin{pmatrix} 0 \\ dx_{123} \end{pmatrix}, \begin{pmatrix} 0 \\ dx_{124} \end{pmatrix}$ $\begin{pmatrix} 0 \\ dx_{134} \end{pmatrix}, \begin{pmatrix} 0 \\ dx_{234} \end{pmatrix}$	$\begin{pmatrix} 0 \\ dx_{1234} \end{pmatrix}$
$H^k(\text{Cone}(c(\omega)))$	$\begin{pmatrix} q_{123} \\ 0 \end{pmatrix}, \begin{pmatrix} q_{124} \\ 0 \end{pmatrix}, \begin{pmatrix} q_{234} \\ 0 \end{pmatrix},$ $\begin{pmatrix} q_{234} \\ 0 \end{pmatrix}, \begin{pmatrix} 0 \\ q_{12} - q_{34} \end{pmatrix}$	$\begin{pmatrix} 0 \\ q_{123} \end{pmatrix}, \begin{pmatrix} 0 \\ q_{124} \end{pmatrix}$ $\begin{pmatrix} 0 \\ q_{134} \end{pmatrix}, \begin{pmatrix} 0 \\ q_{234} \end{pmatrix}$	$\begin{pmatrix} 0 \\ q_{1234} \end{pmatrix}$

Table 1: Cohomology of $\text{Cone}(\omega)$ versus $\text{Cone}(c(\omega))$ on $(T^4, \omega = dx_1 \wedge dx_2 + dx_3 \wedge dx_4)$.

Example 4.4. Let (M, ω) be the six-dimensional, closed, symplectic manifold constructed by Cho in [4] where the symplectic form ω is not hard Lefschetz type. Topologically, M can be described as a two-sphere bundle over a projective $K3$ surface and also has the following properties [4, Theorem 1.3]: (i) M is simply-connected; (ii) the odd degree cohomologies vanish, i.e. $H_{dR}^1(M) = H_{dR}^3(M) = H_{dR}^5(M) = 0$.

Consider the cohomology $PH(M, \omega) \cong H(\text{Cone}(\omega))$. From (4.3), we find

$$\begin{aligned}
b_0^\omega &= b_7^\omega = 1, \\
b_1^\omega &= b_6^\omega = 0, \\
b_2^\omega &= b_5^\omega = b_2(X) - 1, \\
b_3^\omega &= \dim [\ker (\omega : H^2(X) \rightarrow H^4(X))] > 0, \\
b_4^\omega &= \dim [\text{coker} (\omega : H^2(X) \rightarrow H^4(X))] > 0.
\end{aligned}$$

Note that $b_3^\omega = b_4^\omega > 0$ since (M, ω) is not hard Lefschetz, which implies that the map, $\omega : H_{dR}^2(M) \rightarrow H_{dR}^4(M)$, can not be an isomorphism.

For the cone Morse complex and inequalities, we can again choose to work with a

perfect Morse function on M . That such exists is due to a result of Smale [10, Theorem 6.3] which states that any simply-connected manifold of dimension greater than five that has no homology torsion has a perfect Morse function. (No homology torsion here can be seen from applying the Gysin sequence to M as a two-sphere bundle over $K3$.) Since M has trivial odd-degree cohomology, this implies that

$$\begin{aligned} m_0 &= m_6 = 1, \\ m_1 &= m_3 = m_5 = 0, \\ m_2 &= m_4 = b_2(M). \end{aligned}$$

It is straightforward to check that the bounds (4.1)-(4.2) are satisfied. In particular, for the weak cone Morse bound of (4.1), the $k = 3, 4$ case corresponds to

$$\begin{aligned} b_3^\omega &\leq m_3 + m_2 - v_2 = m_2 - v_2, \\ b_4^\omega &\leq m_4 + m_3 - v_2 = m_4 - v_2, \end{aligned}$$

The above demonstrates the necessity of having both the m_k and the m_{k-1} term in the symplectic cone Morse inequalities.

Remark 4.5. We comment that there is a preprint [9] that presents some symplectic Morse-type inequalities which are different from those here and actually not valid generally. For instance, the inequality in [9, Corollary 3] can be expressed in our notation as $\dim F^p H^{n+p+1}(M, \omega) \leq m_{n-p}$, which is not satisfied in the above Cho's non-Kähler six-dimensional example (M, ω) for a perfect Morse function. Specifically, it gives for $p = 0$ case the inequality relation, $b_4^\omega \leq m_3 = 0$, which is inconsistent with $b_4^\omega > 0$ with ω being of non-hard Lefschetz type.

5 Discussion

Thus far, in this paper, we have for simplicity focused on the $p = 0$ case of the TTY cohomologies, $F^p H(M, \omega) = H(\text{Cone}(\omega^{p+1}))$. Let us comment in this final section the $p > 0$ case and lay out the results which generalize the $p = 0$ case. The cone Morse theory in the $p > 0$ case can be considered analytically similar to the computations in this paper though the calculations are more involved. In general, the TTY cohomologies for

$p = 0, 1, \dots, n - 1$ algebraically correspond to [12, 14]:

$$\begin{aligned} F^p H^k(M, \omega) \cong H^k(\text{Cone}(\omega^{p+1})) \cong \text{coker} \left([\omega^{p+1}] : H_{dR}^{k-2p-2}(M) \rightarrow H_{dR}^k(M) \right) \\ \oplus \ker \left([\omega^{p+1}] : H_{dR}^{k-2p-1}(M) \rightarrow H_{dR}^{k+1}(M) \right) \end{aligned} \quad (5.1)$$

with $k = 0, 1, \dots, 2n + 2p + 1$. The relevant cone complex for the general p case would have the elements and differential

$$\text{Cone}^k(\omega^{p+1}) = \Omega^k(M) \oplus \theta \Omega^{k-2p-1}(M), \quad d_C = \begin{pmatrix} d & \omega^{p+1} \\ 0 & -d \end{pmatrix}, \quad (5.2)$$

where θ is now a formal $(2p + 1)$ -form such that $d\theta = \omega^{p+1}$. And the corresponding cone Morse cochain complex would be

$$\text{Cone}^k(c(\omega^{p+1})) = C^k(M, f) \oplus C^{k-2p-1}(M, f), \quad \partial_C = \begin{pmatrix} \partial & c(\omega^{p+1}) \\ 0 & -\partial \end{pmatrix}, \quad (5.3)$$

with $c(\omega^{p+1}) : C^k(M, f) \rightarrow C^{k+2p+2}(M, f)$ given by

$$c(\omega^{p+1}) q_k = \sum_{r_{k+2p+2}} \left(\int_{\overline{\mathcal{M}(r_{k+2p+2}, q_k)}} \omega^{p+1} \right) r_{k+2p+2}, \quad (5.4)$$

which integrates ω^{p+1} over the $2(p+1)$ -dimensional submanifold $\overline{\mathcal{M}(r_{k+2p+2}, q_k)}$ of gradient flow lines from the index $k + 2p + 2$ critical point, r_{k+2p+2} , to q_k .

We can take the inner product on $\text{Cone}^k(\omega^{p+1})$, just as in (2.2), to be

$$\langle \eta_k + \theta \xi_{k-2p-1}, \eta'_k + \theta \xi'_{k-2p-1} \rangle_C = \langle \eta_k, \eta'_k \rangle + \langle \xi_{k-2p-1}, \xi'_{k-2p-1} \rangle. \quad (5.5)$$

This inner product defines the adjoint operator d_C^* that goes into the cone Laplacian $\Delta_C = d_C d_C^* + d_C^* d_C$, which is a second-order elliptic operator on $\text{Cone}^k(\omega^{p+1})$. The Witten deformation method can be applied to this cone Laplacian Δ_C for $p > 0$ following the steps described in Section 2 and 3. The calculations are similar to the $p = 0$ case. At large t , there are again only two generators to the solutions of the deformed harmonic Laplacian localized at each critical point $p \in \text{Crit}(f)$. For instance, for $k \leq n + p$, the

generators take the form

$$\begin{pmatrix} \zeta_k \\ 0 \end{pmatrix}, \quad \begin{pmatrix} -\tau \wedge \omega^p \wedge \zeta_{k-2p-1} \\ \zeta_{k-2p-1} \end{pmatrix},$$

which generalizes the generators in (2.13) and (2.16), respectively. Generalizations of the estimates similar to those for the $p = 0$ case can be carried out which results in the isomorphism of the cohomologies of the cone complex (5.2) with that of the cone Morse complex (5.3). From these results, we can likewise obtain cone Morse inequalities. Explicitly, using the notation

$$s_k^p = \dim F^p H^k(M, \omega) = \dim H^k(\text{Cone}(\omega^{p+1}))$$

to denote the dimension of the TTY cohomology, we expect the following weak and strong cone Morse inequalities for all $p = 0, 1, \dots, n - 1$:

$$s_k^p \leq m_k - v_{k-2p-2} + m_{k-2p-1} - v_{k-2p-1}, \quad (5.6)$$

$$\sum_{i=0}^k (-1)^{k-i} s_i^p \leq \left(\sum_{i=k-2p}^k (-1)^{k-i} m_i \right) - v_{k-2p-1}, \quad (5.7)$$

where $k = 0, 1, \dots, 2n + 2p + 1$ and

$$v_k = \text{rank} \left(c(\omega^{p+1}) : C^k(M, f) \rightarrow C^{k+2p+2}(M, f) \right). \quad (5.8)$$

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