

Stress Functions for Nonlinear Shallow Shell Theory*

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This note obtains the stress function solution of the homogeneous equilibrium equations of the nonlinear shallow shell theory as formulated in [1].

The following is presented as a supplement to the results of the preceding paper [1]. Inspection of the differential equations of equilibrium (1a, b), (5), (33), (34) and (35) in [1], and some mathematical experimentation which makes use of the three linear compatibility equations (12a, b) and (16) in [1], indicate that the six homogeneous nonlinear equilibrium equations may be satisfied identically by means of six stress functions H_1, H_2, F, K_1, K_2 and J , in the form

$$\begin{aligned} N_{11} &= K_{2,2}, & N_{12} &= K_{1,2}, & N_{21} &= -K_{2,1}, & N_{22} &= -K_{1,1}, \\ P_1 &= -K_2 + F_{,2}, & P_2 &= -K_1 - F_{,1}, \end{aligned} \quad (1)$$

$$\begin{aligned} Q_1 &= J_{,2} + (z_{,22} - \kappa_{22})K_1 + (z_{,12} - \kappa_{21})K_2, \\ Q_2 &= -J_{,1} - (z_{,11} - \kappa_{11})K_2 - (z_{,12} - \kappa_{12})K_1, \end{aligned} \quad (2)$$

$$\begin{aligned} M_{11} &= H_{2,2} - (z_{,22} - \kappa_{22})F + \gamma_2 K_2, \\ M_{22} &= -H_{1,1} - (z_{,11} - \kappa_{11})F - \gamma_1 K_1, \\ M_{12} &= H_{1,2} + (z_{,12} - \kappa_{21})F + \gamma_2 K_1 - J, \\ M_{21} &= -H_{2,1} + (z_{,12} - \kappa_{12})F - \gamma_1 K_2 + J. \end{aligned} \quad (3)$$

Differential equations for the determination of the six stress functions follow upon introducing strain measures in terms of stress functions into the equations of compatibility. In order to express the κ_{ij} and γ_i in terms of stress functions, it is necessary to make use of the relevant stress strain relations. To illustrate, for stress strain relations of the form

$$\begin{aligned} M_{11} &= D(\kappa_{11} + \nu\kappa_{22}), & M_{22} &= \dots, & M_{12} &= D_S\kappa_{12}, & M_{21} &= \dots, \\ \gamma_i &= A_Q Q_i, \end{aligned} \quad (4)$$

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the six strain measures κ_{ij} and γ_i are determined in terms of stress functions upon solving a system of six simultaneous linear equations, obtained by substituting (4) into the left hand sides of (2) and (3),

$$\begin{aligned} \gamma_1 + A_Q K_1 \kappa_{22} + A_Q K_2 \kappa_{21} &= A_Q (J_{,2} + z_{,22} K_1 + z_{,12} K_2), \\ \gamma_2 - A_Q K_2 \kappa_{11} - A_Q K_1 \kappa_{12} &= -A_Q (J_{,1} + z_{,11} K_2 + z_{,12} K_1). \end{aligned} \quad (5)$$

$$\begin{aligned} -K_2 \gamma_2 + D \kappa_{11} + (vD - F) \kappa_{22} &= H_{2,2} - z_{,22} F, \\ K_1 \gamma_1 + (vD - F) \kappa_{11} + D \kappa_{22} &= -H_{1,1} - z_{,11} F. \end{aligned} \quad (6)$$

$$\begin{aligned} -K_1 \gamma_2 + D_S \kappa_{12} + F \kappa_{21} &= H_{1,2} + z_{,12} F - J, \\ K_2 \gamma_1 + F \kappa_{12} + D_S \kappa_{21} &= -H_{2,1} + z_{,12} F + J. \end{aligned} \quad (7)$$

We limit ourselves here to the solution of (5) to (7) for the case $A_Q = 0$, for which (5) becomes $\gamma_i = 0$, and then from (6)

$$\begin{aligned} \kappa_{11} &= \frac{DH_{2,2} + (vD - F)H_{1,1} + [z_{,11}(vD - F) - z_{,22}D]}{(1 - v^2)D^2 + 2vDF - F^2}, \\ \kappa_{22} &= -\frac{DH_{1,1} + (vD - F)H_{2,2} - [z_{,22}(vD - F) - z_{,11}D]F}{(1 - v^2)D^2 + 2vDF - F^2}. \end{aligned} \quad (8)$$

In solving (7), we observe that equation (16) of [1], with $\gamma_i = 0$, implies that $\kappa_{12} = \kappa_{21}$. Therewith,

$$\kappa_{12} = \kappa_{21} = \frac{H_{1,2} - H_{2,1} + 2z_{,12}F}{2(D_S + F)}, \quad J = \frac{H_{1,2} + H_{2,1}}{2}. \quad (9)$$

With (8) and (9) the stress strain relations (4) give as expressions for stress couples

$$\begin{aligned} \frac{M_{12}}{D_S} &= \frac{M_{21}}{D_S} = \frac{2z_{,12}F + H_{1,2} - H_{2,1}}{2(D_S + F)}, \\ \frac{M_{11}}{D} &= \frac{[(1 - v^2)D + vF](H_{2,2} - z_{,22}F) - F(H_{1,1} + z_{,11}F)}{(1 - v^2)D^2 + 2vDF - F^2}, \\ \frac{M_{22}}{D} &= -\frac{[(1 - v^2)D + vF](H_{1,1} + z_{,11}F) - F(H_{2,2} - z_{,22}F)}{(1 - v^2)D^2 + 2vDF - F^2}, \end{aligned} \quad (10)$$

and the equations (2) give as expressions for the transverse shear resultants

$$\begin{aligned} Q_1 &= J_{,2} + K_2 \frac{2z_{,12}D_S + H_{2,1} - H_{1,2}}{2(D_S + F)} \\ &\quad + K_1 \frac{z_{,22}D[(1 - v^2)D + vF] + z_{,11}DF + DH_{1,1} + (vD - F)H_{2,2}}{(1 - v^2)D^2 + 2vDF - F^2}, \\ Q_2 &= -J_{,1} - K_1 \frac{2z_{,12}D_S + H_{2,1} - H_{1,2}}{2(D_S + F)} \\ &\quad - K_2 \frac{z_{,11}D[(1 - v^2)D + vF] + z_{,22}DF + DH_{2,2} - (vD - F)H_{1,1}}{(1 - v^2)D^2 + 2vD - F^2}. \end{aligned} \quad (11)$$

Reference

1. E. REISSNER, "On the equations of nonlinear shallow shell theory", *Studies in Applied Mathematics* **48**, 2, 171–175, (1969).

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